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**Electromagnetic Fields
of Magnetized Compact
Stars in Braneworld**

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Plan of talk



- Basic assumptions and approach
- Electromagnetic Fields of Magnetized Star in Brane
- Stationary Solutions to Maxwell Equations for Magnetized Highly Conducting Spherical Star in Braneworld
- Astrophysical Consequences
- Conclusion

Assumptions and Approach

- GR effects are important for compact objects

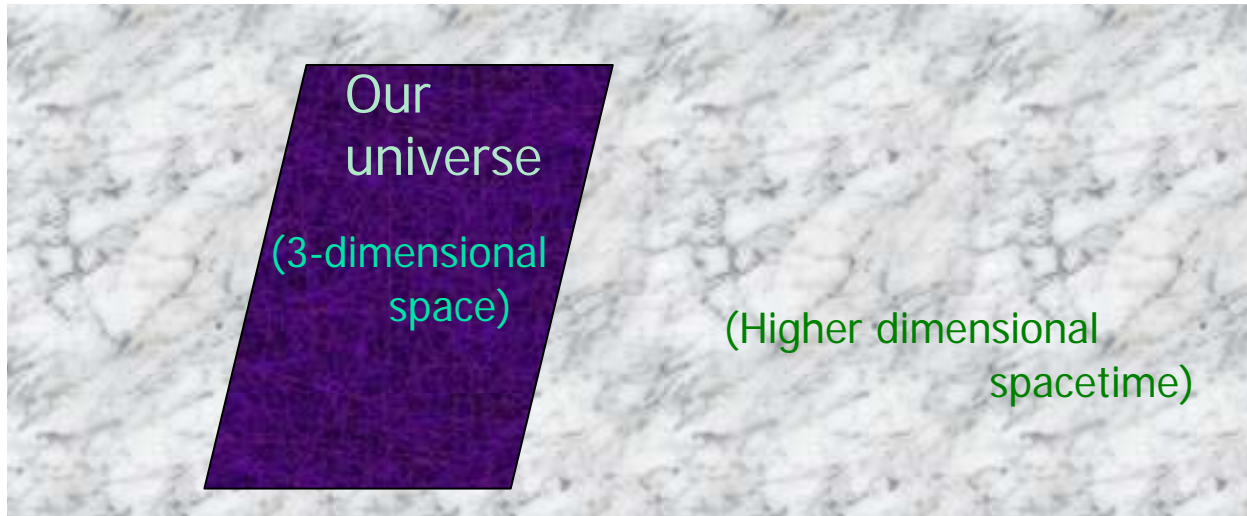
$$R \sim 10 \text{ km}, R_g = \frac{2GM}{c^2 r} \sim 5 \text{ km}, \left(1 - \frac{2GM}{c^2 r}\right) \sim 0.5$$

- Immense difficulty of simultaneously solving the Maxwell eqs and the highly nonlinear Einstein eqs
- Perfect conductivity in the interior region of the star
- The exterior is considered to be as vacuum
- Gravitational field feedback is neglected

■ Braneworld

Recent progress of superstring theory says that our four dimensional universe is entirely restricted to a brane inside a higher dimensional space, called the bulk. According to RS model, the universe is 5D AdS space and elementary particles except of graviton is localized at 3+1 spacetime.

The new picture of the universe



$$ds^2 = -A^2(r)dt^2 + H^2(r)dr^2 + r^2 d\theta^2 + r^2 \sin^2 \theta d\phi^2$$

EINSTEIN EQS FOR UNKNOWN METRIC FUNCTIONS A (r) & H(r)

$$\frac{1}{r^2} - \frac{1}{H^2} \left(\frac{1}{r^2} - \frac{2}{r} \frac{H'}{H} \right) = 8\pi \rho^{eff} ,$$

$$-\frac{1}{r^2} + \frac{1}{H^2} \left(\frac{1}{r^2} + \frac{2}{r} \frac{A'}{A} \right) = 8\pi \left(\rho^{eff} + \frac{4}{k^2 \lambda} P \right) ,$$

$$p' + \frac{A'}{A} (\rho + p) = 0 ,$$

$$U' + 4 \frac{A'}{A} U + 2P' + 2 \frac{A'}{A} P + \frac{6}{r} P = -2(4\pi)^2 (\rho + p) \rho' .$$

nonlocal energy density U and nonlocal pressure P .

ρ^{eff} is the effective total energy density

ρ and p are matter energy density and pressure

OUTSIDE THE STAR

$$A^2(r) \equiv \left(1 - \frac{2M}{r} + \frac{Q}{r^2} \right) = H^{-2}(r) , \quad r > R$$

negative Weyl “charge” $Q = -3MR\rho/\lambda$

M and R are the total mass

and radius of the star, λ is the brane tension.

HORIZON

$$r_+ = M \left(1 + \sqrt{1 - \frac{Q^2}{M^2}} \right)$$

MAXWELL EQS

$$e^{\alpha\beta\mu\nu} F_{\beta\mu}{}_{,\nu} = 0,$$

$$F^{\alpha\beta}{}_{;\beta} = 4\pi J^\alpha$$

FIRST PAIR

$$\sin \theta (r^2 B^{\hat{r}})_{,r} + Hr (\sin \theta B^{\hat{\theta}})_{,\theta} + Hr B^{\hat{\phi}}_{,\phi} = 0 ,$$

$$(r \sin \theta) \frac{\partial B^{\hat{r}}}{\partial t} = A \left[E^{\hat{\theta}}_{,\phi} - (\sin \theta E^{\hat{\phi}})_{,\theta} \right] ,$$

$$(Hr \sin \theta) \frac{\partial B^{\hat{\theta}}}{\partial t} = -AHE^{\hat{r}}_{,\phi} + \sin \theta (rAE^{\hat{\phi}})_{,r} ,$$

$$(Hr) \frac{\partial B^{\hat{\phi}}}{\partial t} = - (rAE^{\hat{\theta}})_{,r} + AHE^{\hat{r}}_{,\theta}$$

SECOND PAIR

$$\sin \theta \left(r^2 E^{\hat{r}} \right)_{,r} + Hr \left(\sin \theta E^{\hat{\theta}} \right)_{,\theta} + Hr E^{\hat{\phi}}_{,\phi} = 4\pi Hr^2 \sin \theta J^{\hat{t}} ,$$

$$A \left[\left(\sin \theta B^{\hat{\phi}} \right)_{,\theta} - B^{\hat{\theta}}_{,\phi} \right] = (r \sin \theta) \frac{\partial E^{\hat{r}}}{\partial t} + 4\pi Ar \sin \theta J^{\hat{r}} ,$$

$$AH B^{\hat{r}}_{,\phi} - \sin \theta \left(r AB^{\hat{\phi}} \right)_{,r} = (Hr \sin \theta) \frac{\partial E^{\hat{\theta}}}{\partial t} + 4\pi AHr \sin \theta J^{\hat{\theta}} ,$$

$$\left(Ar B^{\hat{\theta}} \right)_{,r} - AH B^{\hat{r}}_{,\theta} = (Hr) \frac{\partial E^{\hat{\phi}}}{\partial t} + 4\pi AHr J^{\hat{\phi}} .$$

Dipolar Magnetic Field and Separable Solutions

$$B^{\hat{r}}(r, \theta) = F(r) \cos \theta,$$

$$B^{\hat{\theta}}(r, \theta) = G(r) \sin \theta,$$

$$B^{\hat{\phi}}(r, \theta) = 0.$$



MAXWELL EQS FOR RADIAL FUNCTIONS

$$(r^2 F)_{,r} + 2HrG = 0 ,$$

$$(rAG)_{,r} + AHF = 0 .$$



Interior Analytical Solution

$$F = \frac{C_1}{R^3} \mu,$$

$$G = -\frac{H^{-1}C_1}{R^3} \mu = -H^{-1}F$$

For Stiff Matter

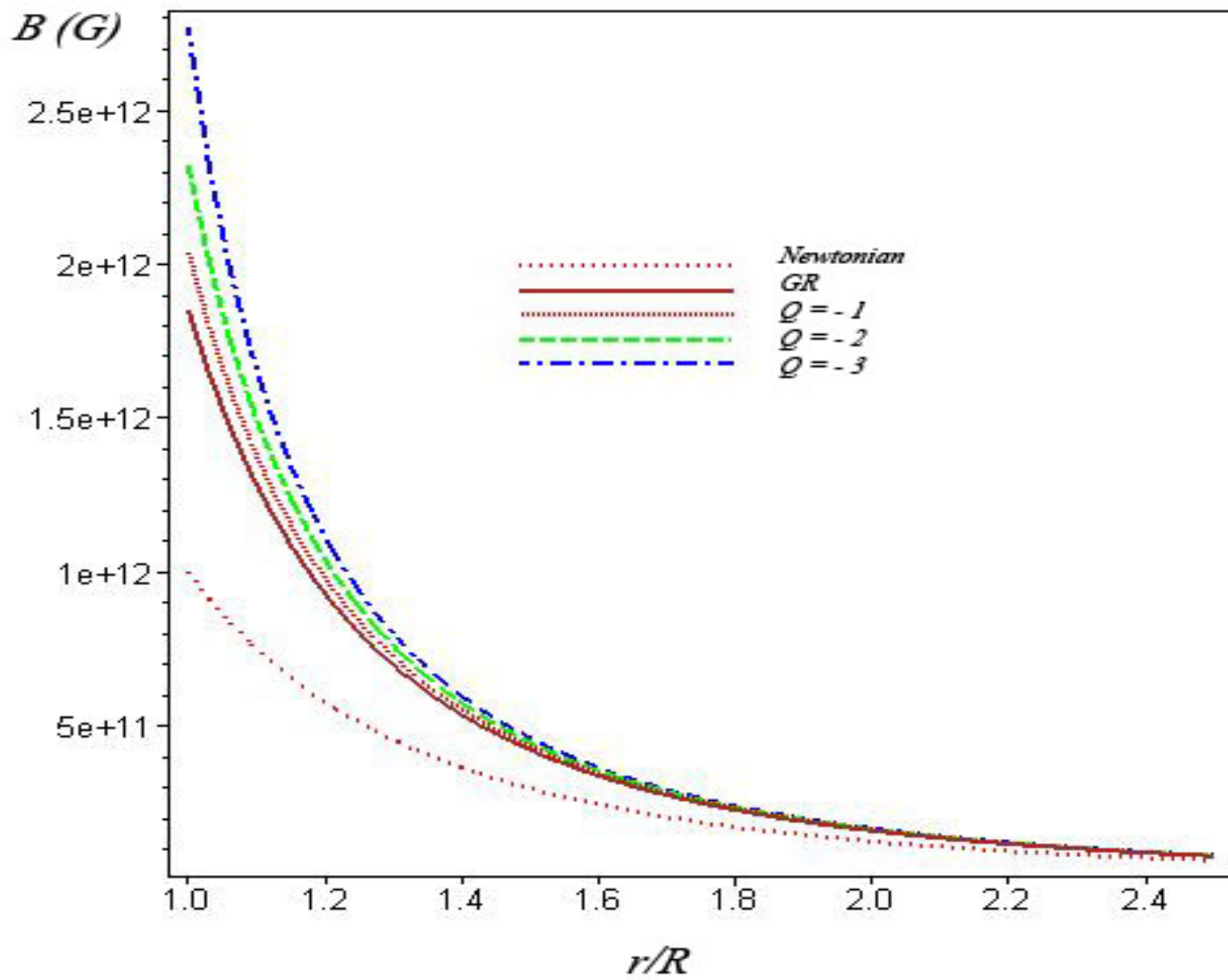
$$\rho^{eff} = p^{eff} + \frac{4}{k^4 \lambda} P$$

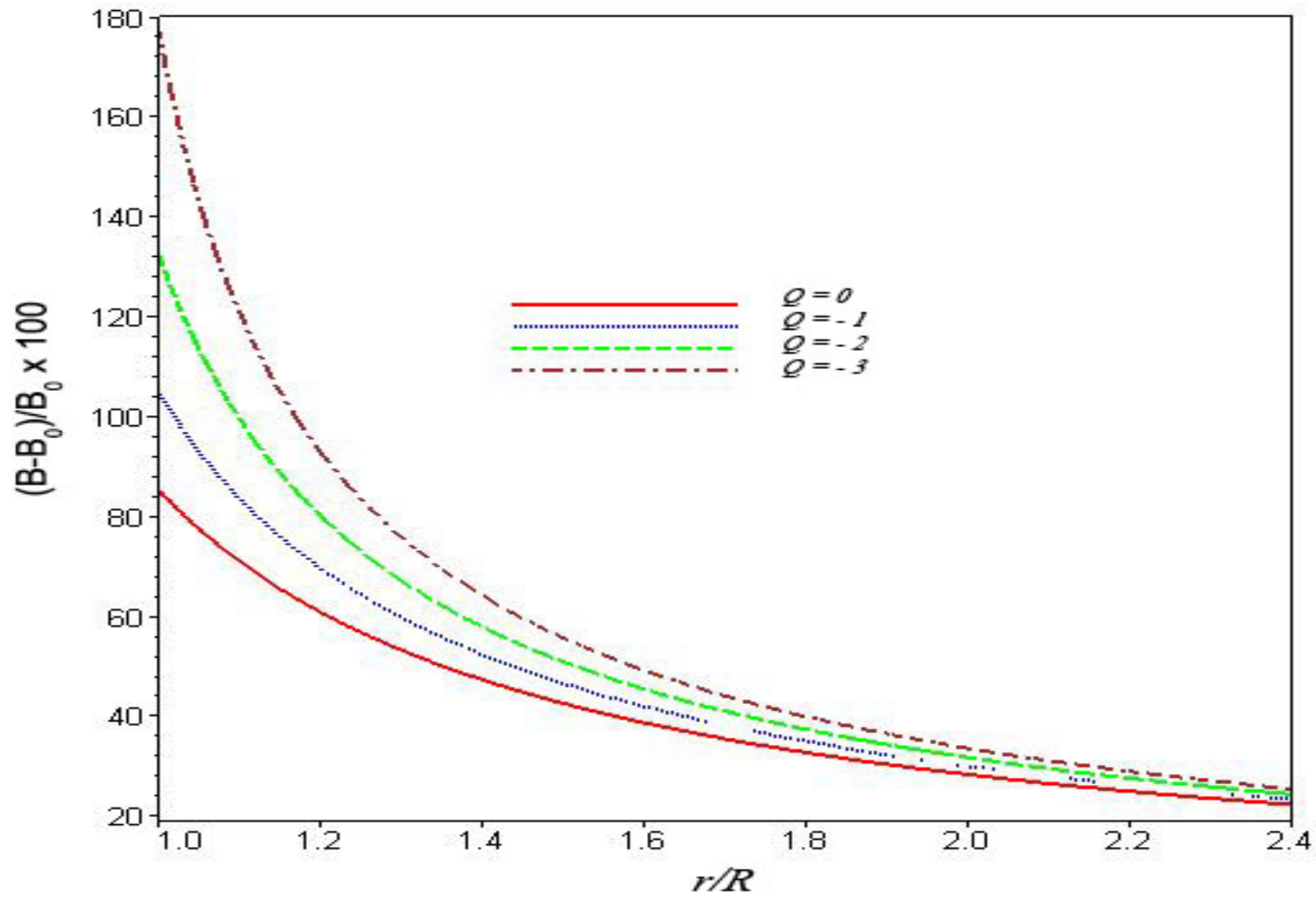
Exterior Numerical Solution for Magnetic Field

$$\frac{d}{dr} \left[\left(1 - \frac{2M}{r} + \frac{Q}{r^2} \right) \frac{d}{dr} (r^2 F) \right] - 2F = 0$$

The analytical solution exists, when $Q=0$:

$$\frac{F_{GR}(r)}{F_{Newt}(r)} = -\frac{3R^3}{8M^3} \left[\ln N^2 + \frac{2M}{r} \left(1 + \frac{M}{r} \right) \right]$$







Interior Numerical Solution

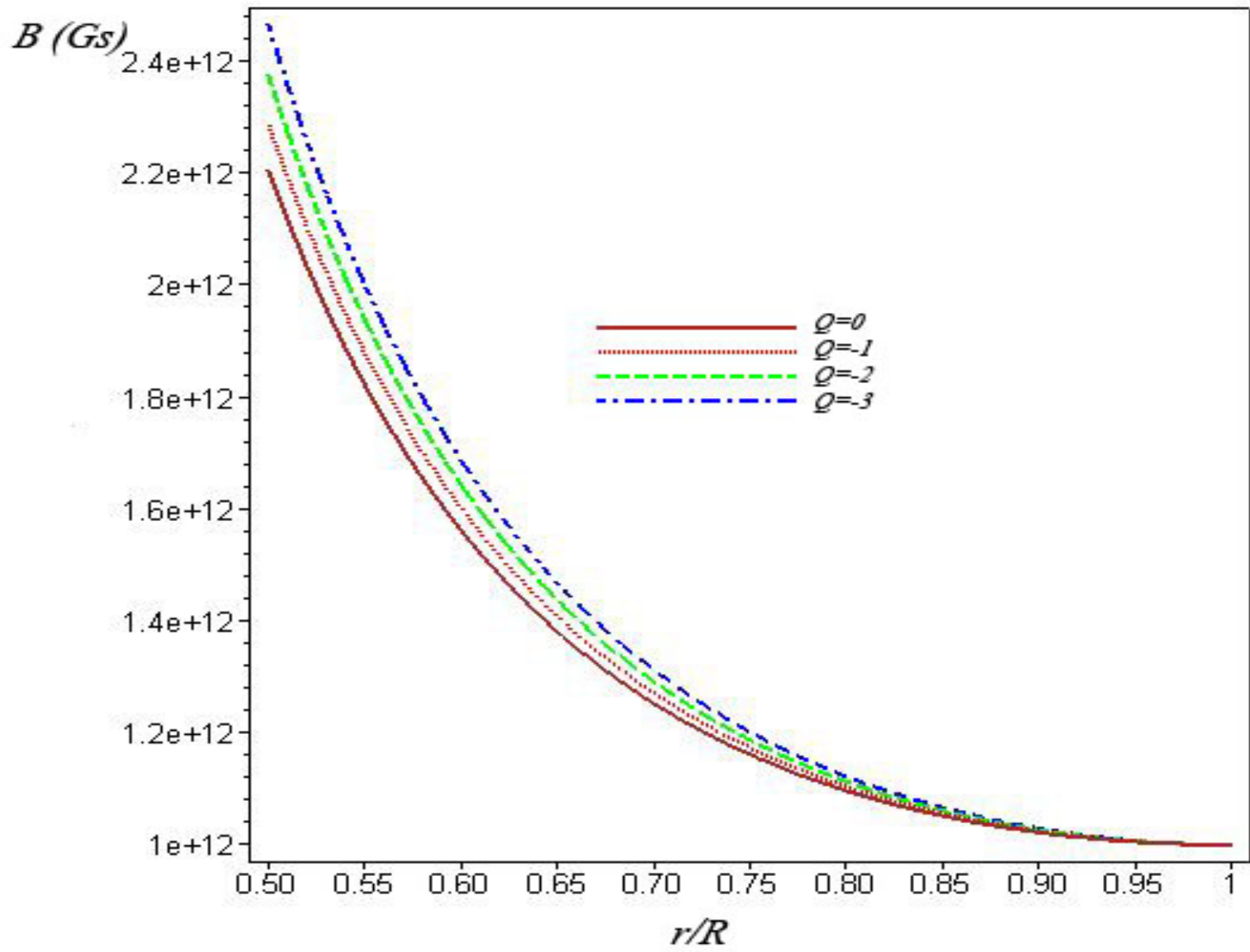
$$\frac{d}{dr} \left[\frac{\Delta(r)}{\rho + p(r)} \frac{d}{dr} (r^2 F) \right] - \frac{2F}{[\rho + p(r)]\Delta(r)} = 0$$

where

$$\Delta(r) \equiv \left[1 - \frac{2M}{r} \left(\frac{r}{R} \right)^3 \left[1 - \frac{Q}{6MR} \right] \right]^{1/2},$$

$$p(r) = \frac{[\Delta(r) - \Delta(R)](1 - Q/3MR)\rho}{[3\Delta(R) - \Delta(r)] - [3\Delta(R) - 2\Delta(r)]Q/3MR},$$

$$m(r) = M \left[1 - \frac{Q}{6MR} \right] \left(\frac{r}{R} \right)^3.$$



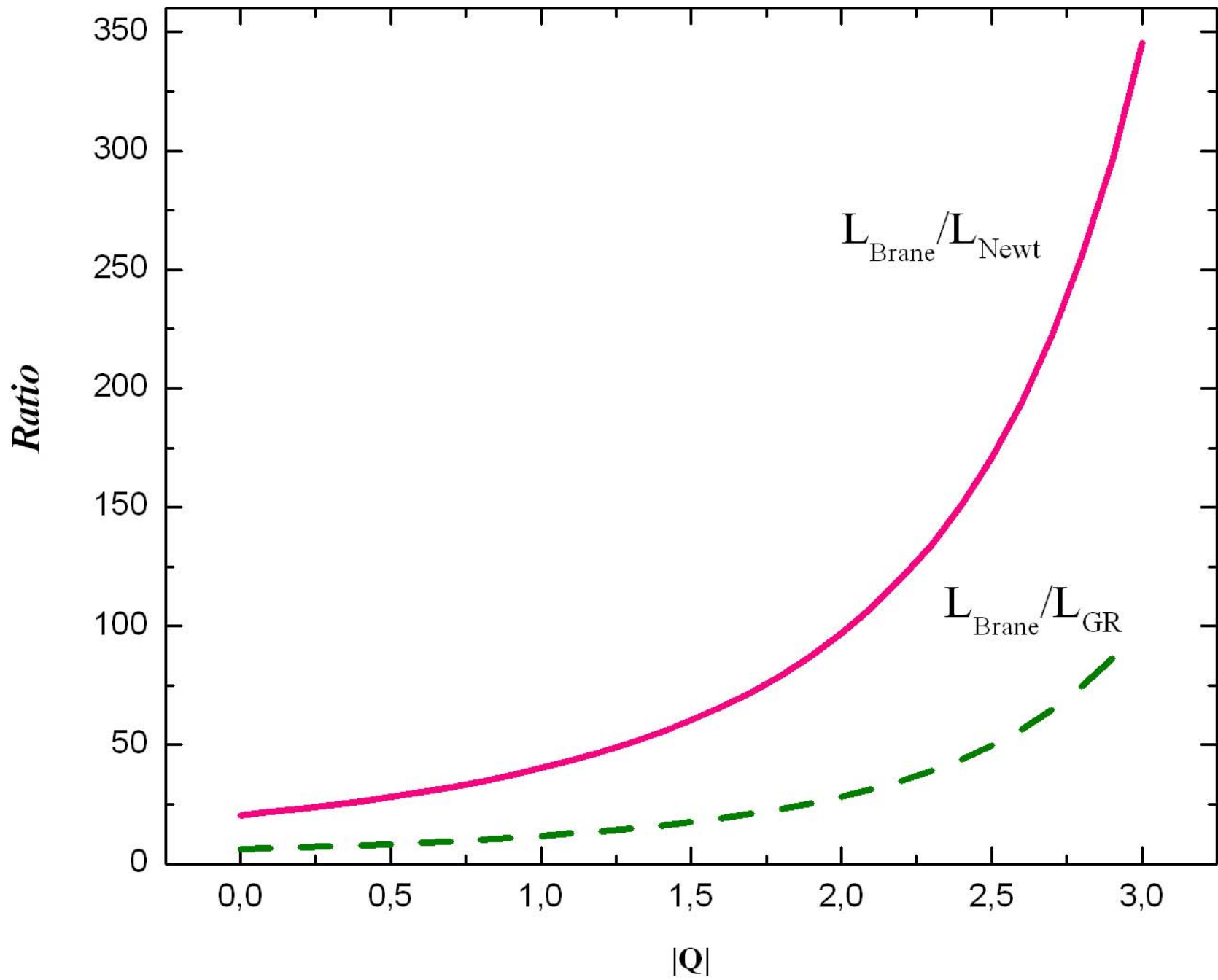
Astrophysical Consequences

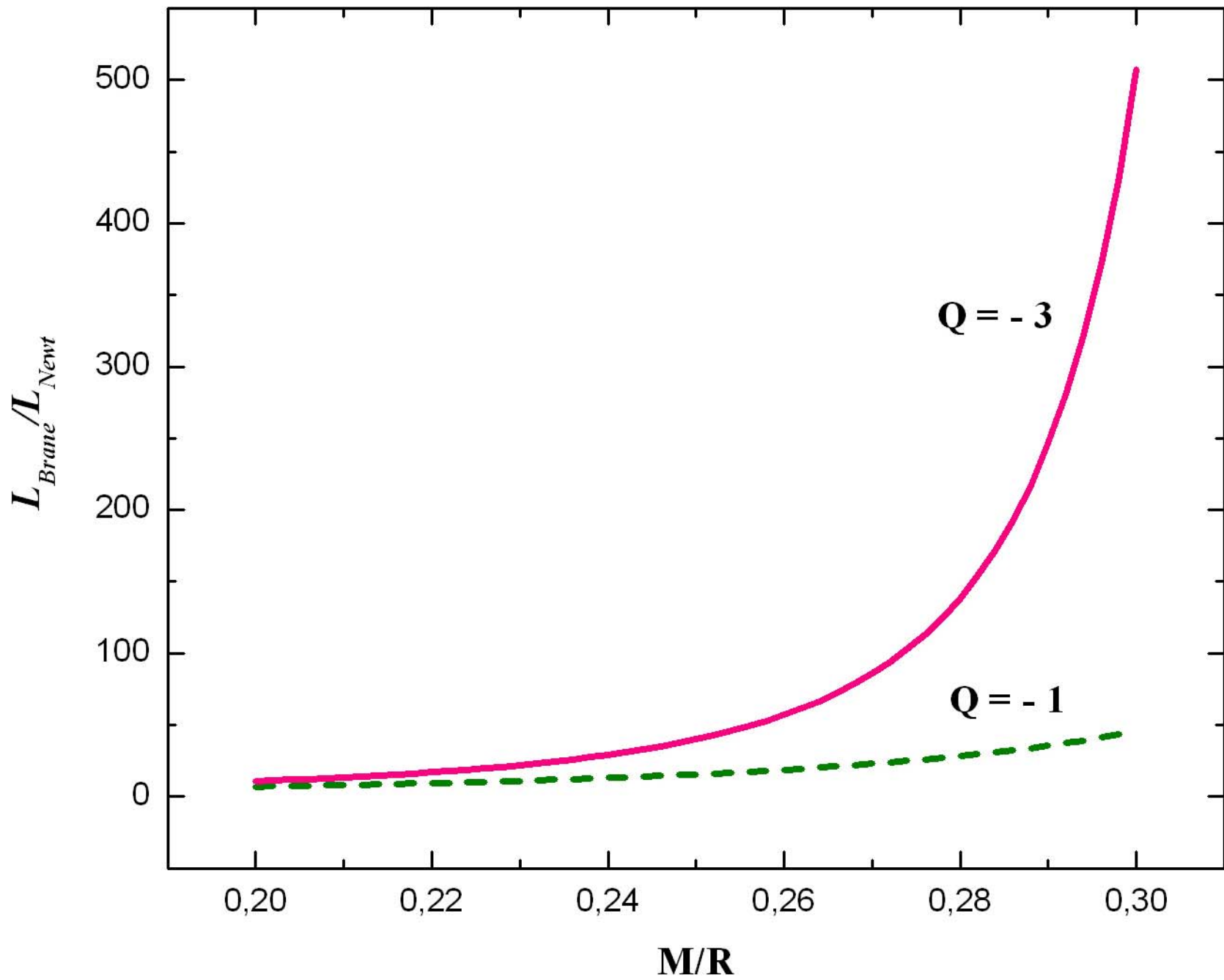
Luminosity of the relativistic star is

$$L_{em} = \frac{\Omega_R^4 R^6 \tilde{B}_0^2}{6c^3} \sin^2 \chi$$

The rate of energy loss is

$$\frac{L_{em}}{(L_{em})_{Newt}} = \left(\frac{F_R}{A_R^2} \right)^2$$





CONCLUSION

- The effect of additional braneworld tension on electromagnetic fields is considered.
- The enhancement of magnetic field coming from the effect of brane tension depends on the parameter Q : For larger Q , the effect is stronger.
- There are two effects of brane tension on magneto-dipolar radiation: a) Due to the amplification of surface magnetic field by brane tension; b) Due to the presence of Q/r^2 , in the red shift factor.
- The expression for magnetodipolar luminosity of rotating braneworld magnetized star gives enhancement up to two orders