



Poster Presentation

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External Electromagnetic Fields of Slowly Rotating Relativistic Magnetized NUT Stars

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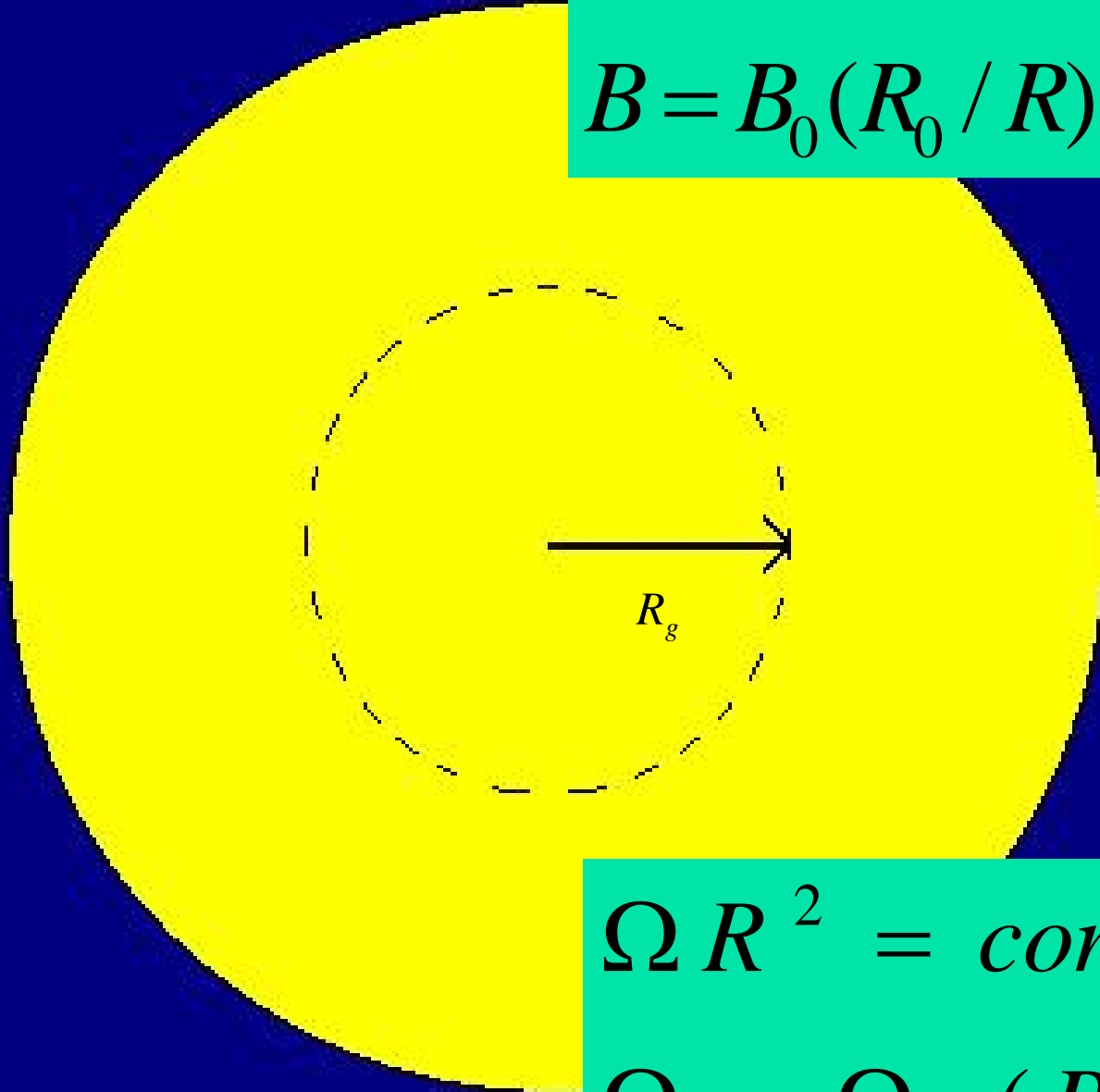
Plan of talk



- **Basic assumptions and approach**
- **Electromagnetic Fields of Slowly Rotating Axial Symmetric Magnetized NUT Star**
- **Conclusion**

$$BR^2 = \text{const} \Rightarrow$$

$$B = B_0 (R_0 / R)^2 \propto 10^{12} G$$



$$\Omega R^2 = \text{const} \Rightarrow$$

$$\Omega = \Omega_0 (R_0 / R)^2$$

$$R \sim 10\text{km}, R_g = \frac{2GM}{c^2 r} \sim 5\text{km}, \left(1 - \frac{2GM}{c^2 r}\right) \sim 0.5$$

- GR effects are important for compact objects
- Immense difficulty of simultaneously solving the Maxwell eqs and the highly nonlinear Einstein eqs

$$e^{\alpha\beta\mu\nu} F_{\beta\mu},{}_{,\nu} = 0,$$

$$F^{\alpha\beta}{}_{;\beta} = 4\pi J^{\alpha}$$

$$R_{\alpha\beta} - \frac{1}{2} g_{\alpha\beta} R = \kappa T_{\alpha\beta},$$

$$T_{\alpha\beta} = T_{(em)\alpha\beta} + T_{(G)\alpha\beta}$$

APPROACH

EMF are considered in a given background geometry

$$R_{\alpha\beta} - g_{\alpha\beta} R/2 = \kappa T_{\alpha\beta}$$

$$T^{\alpha\beta} = T^{\alpha\beta}_{(\text{matter})} + T^{\alpha\beta}_{(\text{em})} \quad \frac{1}{4} T^{\alpha\beta}_{(\text{matter})}$$

$$\rho_{(\text{em})} c^2 \quad \frac{1}{4} B^2/8\pi \quad \rho_{(\text{em})} \quad \frac{1}{4} 10^3 \text{ g/cm}^3 \quad \text{if } B \text{ is } 10^{12} \text{ G}$$

$$\rho_{(\text{em})} \dot{\rho}_{(\text{matter})} \quad \frac{1}{4} 10^{14} \text{ g/cm}^3 \quad \text{consequently } T^{\alpha\beta}_{(\text{em})} \dot{T}^{\alpha\beta}_{(\text{matter})}$$

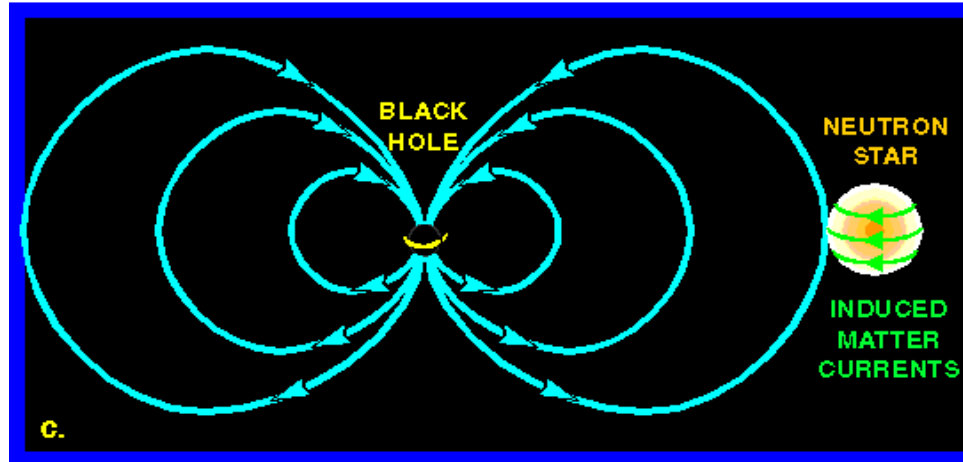
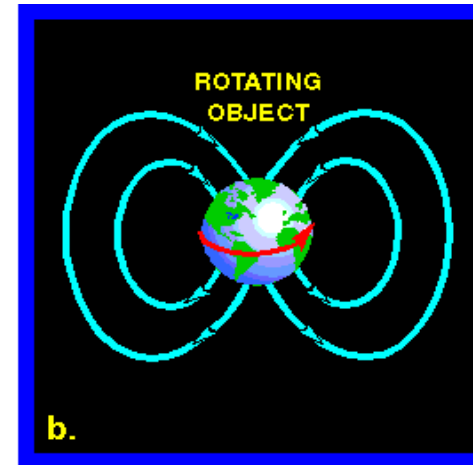
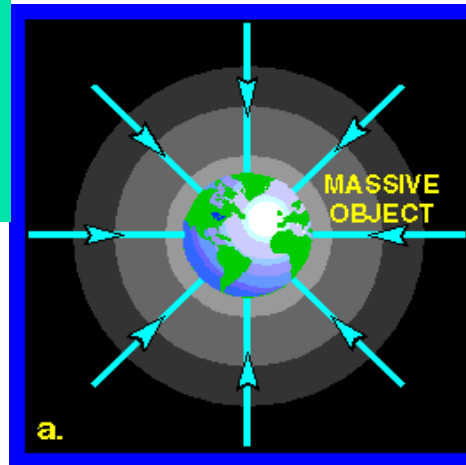
Slow rotation approximation: comparing the neglected terms with the lower order one gives ratios

$$R^3 \Omega^2 / GM < 10\%$$

even for fastest rotating stars (PSR J1748-2446ad & PSR 1937+214)

$$m \frac{d\vec{v}}{dt} = m \left(\vec{E}_g + \frac{1}{c} \vec{v} \times \vec{B}_g \right), \quad \vec{E}_g = -\frac{GM}{r^2} \vec{r}$$

$$\vec{B}_g = \frac{2G}{c} \left[\frac{\vec{J} - 3(\vec{J} \cdot \vec{r})\vec{r}}{r^3} \right]$$



METRIC OF SLOWLY ROTATING NUT STAR

$$ds^2 = -N^2 dt^2 + N^{-2} dr^2 + r^2 d\theta^2 + r^2 \sin^2 \theta d\phi^2 - 2 \left[\omega(r) r^2 \sin^2 \theta + 2lN^2 \cos \theta \right] dt d\phi$$

$$\frac{B^2}{8\pi\mu c^2} \simeq 1.6 \times 10^{-6} \left(\frac{B}{10^{15} \text{ G}} \right)^2 \left(\frac{1.4M_{\odot}}{M} \right) \left(\frac{R}{15 \text{ km}} \right)^3$$

$$\frac{E^2}{8\pi\mu c^2} \simeq 0.2 \times 10^{-20} \left(\frac{E}{3 \times 10^{10} \text{ V/cm}} \right)^2 \left(\frac{1.4M_{\odot}}{M} \right) \left(\frac{R}{15 \text{ km}} \right)^3$$


$$\omega(r) \sim 2aM/r^3$$

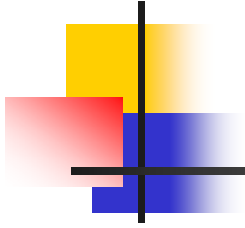
ZAMO OBSERVER:

$$(u^\alpha)_{\text{ZAMO}} \equiv N^{-1} \left(1, 0, 0, \omega + \frac{2l \cos \theta}{r^2 \sin^2 \theta} N^2 \right);$$

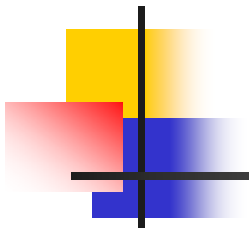
$$(u_\alpha)_{\text{ZAMO}} \equiv N \left(-1, 0, 0, 0 \right).$$

$$N \equiv \left(1 - \frac{2M}{r} \right)^{1/2}$$

Monopolar configuration for MF

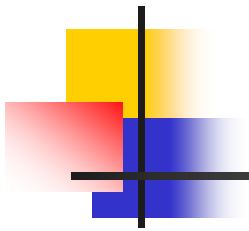


$$B^r = \frac{\mu_0}{r^2}$$



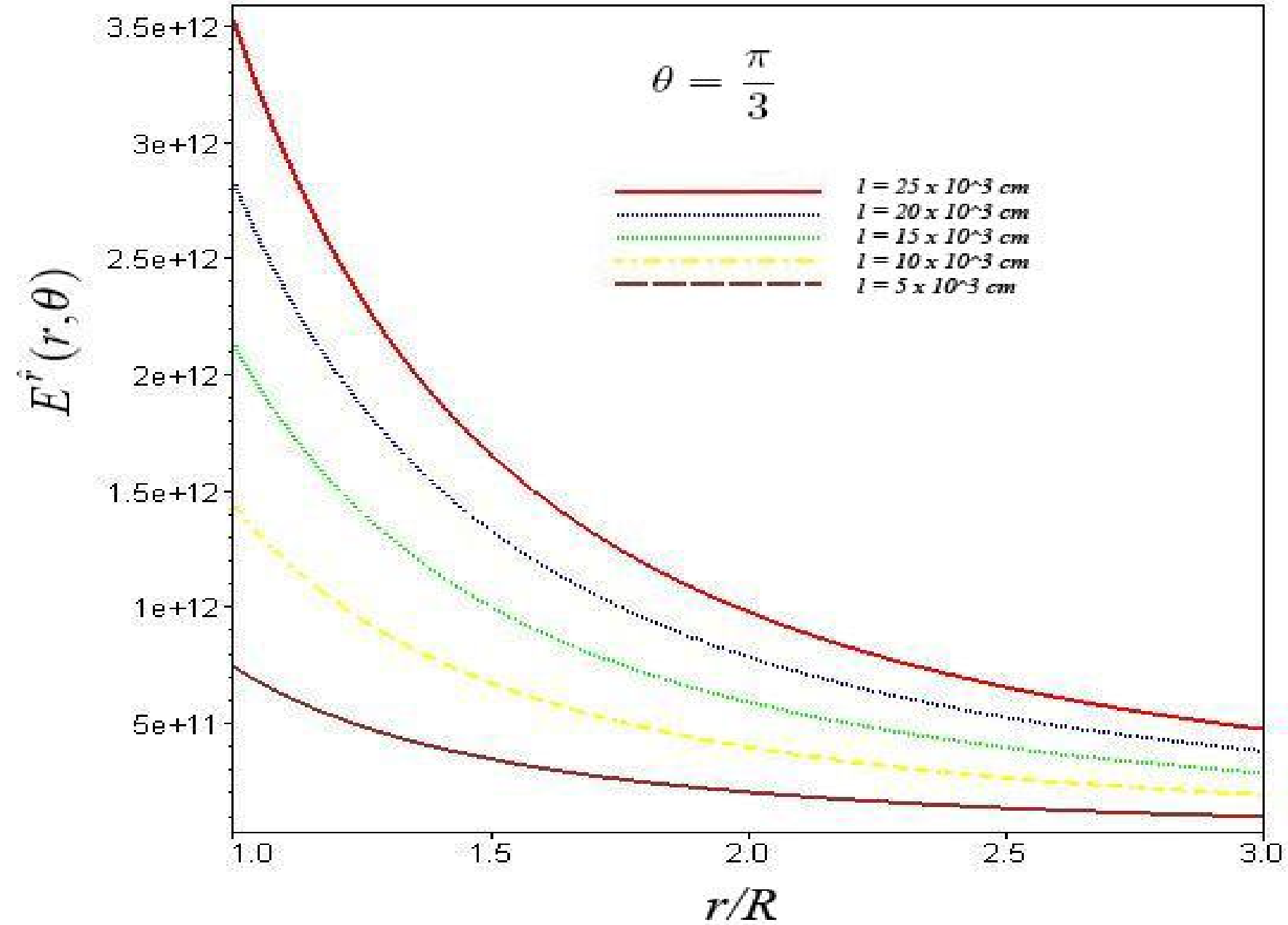
$$E_{,\theta}^{\hat{r}} - (rNE^{\hat{\theta}})_{,r} - \frac{3\omega\mu_0}{r} \sin\theta - \frac{4\mu_0 l}{r^3} \operatorname{ctg}\theta = 0$$

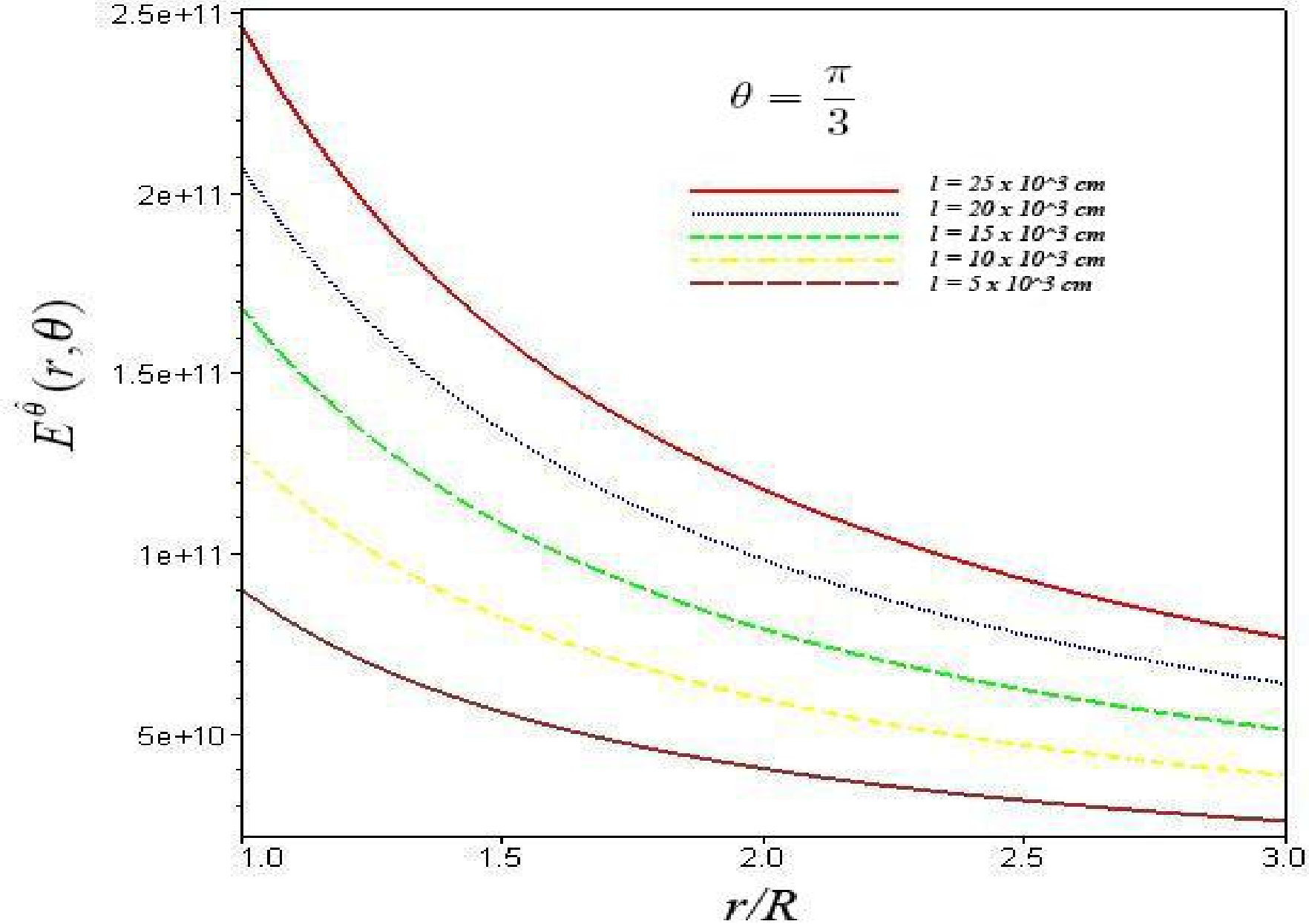
$$\sin\theta (r^2 E^{\hat{r}})_{,r} + \frac{r}{N} (\sin\theta E^{\hat{\theta}})_{,\theta} = 0$$



$$E^{\hat{r}}(r, \theta) = \frac{\Omega R \mu_0}{r^2} - \left(\frac{\omega \mu_0}{M} \right) \cos \theta - \left(\frac{\mu_0 l}{2M^2} \right) \frac{1}{r} \left[1 + \frac{2M}{r} \ln \frac{r}{2M} \right]$$

$$E^{\hat{\theta}}(r, \theta) = -N \frac{\mu_0}{2M} \left\{ \omega \sin \theta + \frac{l}{2Mr} \left(1 + \frac{2M}{r} \right) \operatorname{ctg} \theta \right\}$$

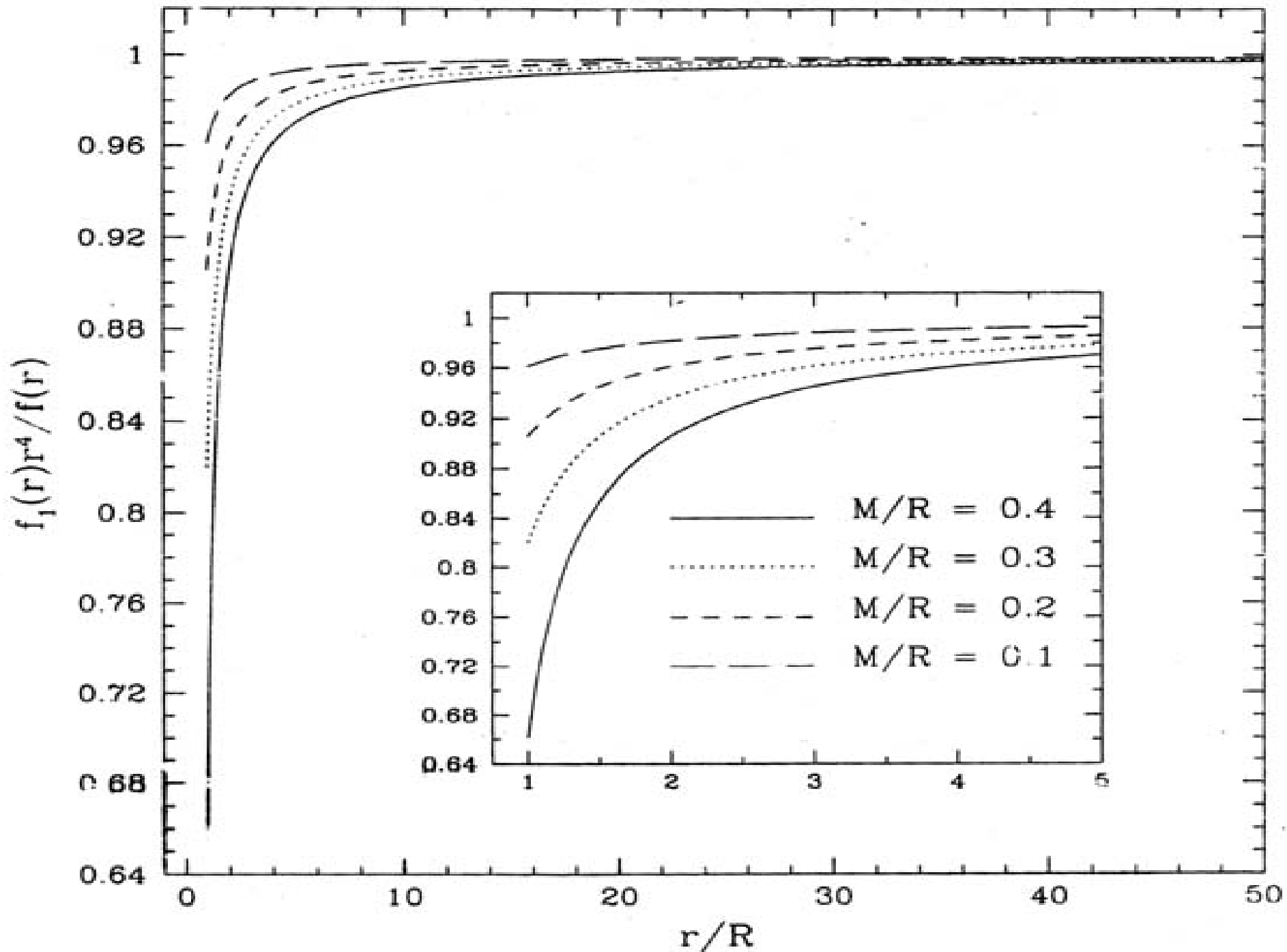


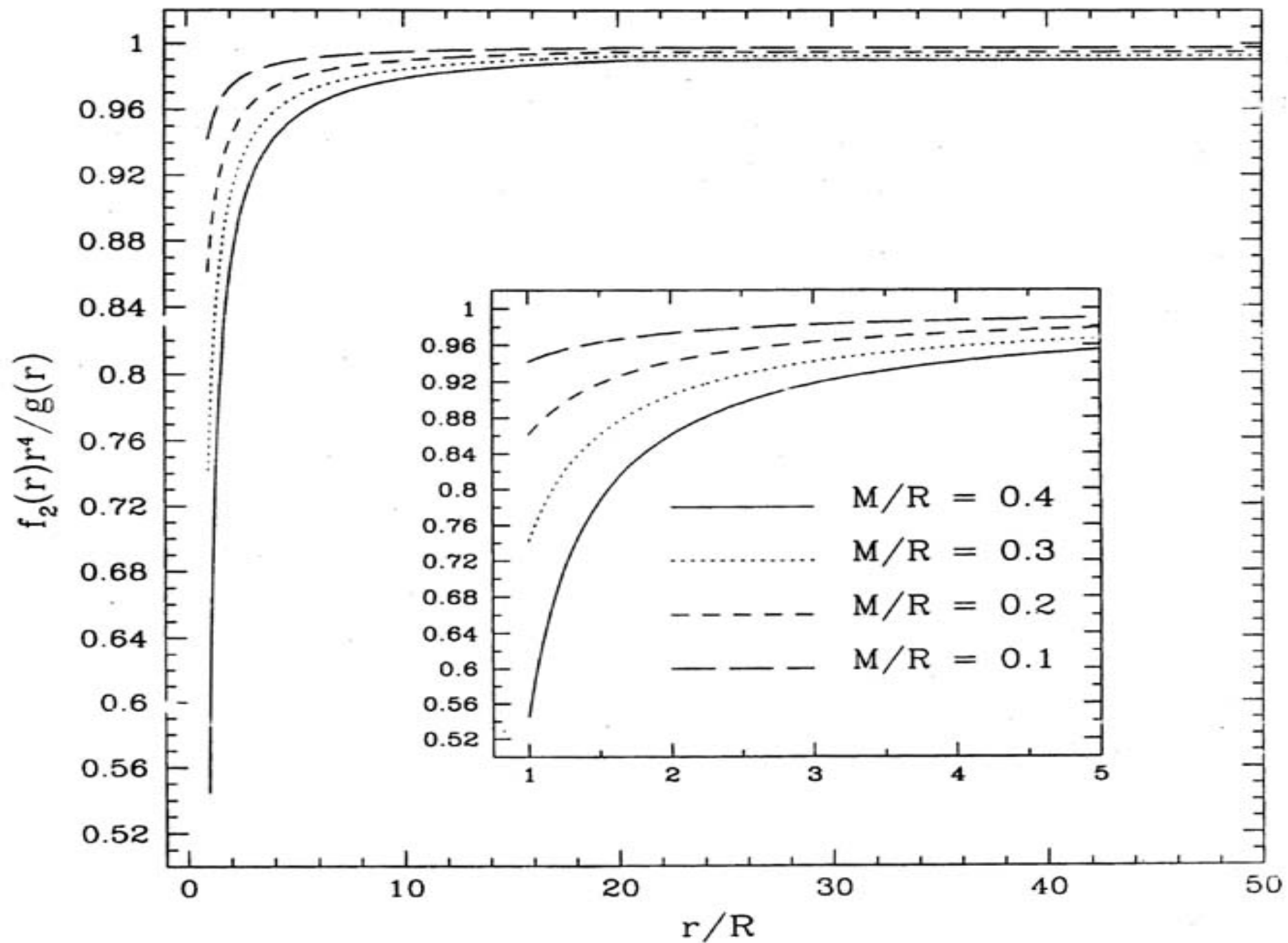


DIPOLAR MAGNETIC FIELD

$$B^r = -\frac{3\mu}{4M^3} \left[\ln N^2 + \frac{2M}{r} \left(1 + \frac{M}{r} \right) \right] \cos \theta,$$

$$B^\theta = \frac{3\mu N}{4M^2 r} \left[\frac{r}{M} \ln N^2 + \frac{1}{N^2} + 1 \right] \sin \theta.$$





$$E^{\hat{r}} = \left\{ \frac{15\omega r^3}{16M^5 c} \left\{ C_3 \left[\left(3 - \frac{2r}{M} \right) \ln N^2 + \frac{2M^2}{3r^2} + \frac{2M}{r} - 4 \right] + \frac{2M^2}{5r^2} \ln N^2 + \frac{4M^3}{5r^3} \right\} \right. \\ \left. + \frac{\Omega}{6cR^2} C_1 C_2 \left[\left(3 - \frac{2r}{M} \right) \ln N^2 + \frac{2M^2}{3r^2} + \frac{2M}{r} - 4 \right] \right\} (3 \cos^2 \theta - 1) \mu - \frac{2\mu l}{3Mr^3} \cos \theta ,$$

$$E^{\hat{\theta}} = - \left\{ \frac{45\omega r^3}{16M^5 c} N \left\{ C_3 \left[\left(1 - \frac{r}{M} \right) \ln N^2 - 2 - \frac{2M^2}{3r^2 N^2} \right] + \frac{4M^4}{15r^4 N^2} \right\} \right. \\ \left. + \frac{\Omega}{2cR^2} C_1 C_2 N \left[\left(1 - \frac{r}{M} \right) \ln N^2 - 2 - \frac{2M^2}{3r^2 N^2} \right] \right\} [2 \sin \theta \cos \theta] \mu - \frac{\mu l}{3Mr^3} \sin \theta .$$

CONCLUSION

- E/M Fields of Rotating Magnetized NUT stars in General Relativity:

GR Effects Resulting from Dragging of

Inertial Frames $\omega = 2 aM / r^3$

Monopolar Part $GM/c^2 R$

and NUT parameter l

are Very Important for EF

For MF only mass of star plays a role