

# ACCRETING ISOLATED NEUTRON STARS

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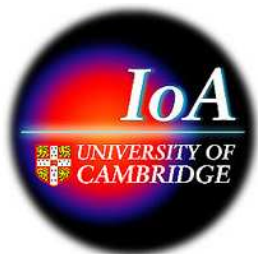
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- “Missing” Accreting Isolated Neutron Stars
- Does this challenge the Theory of Accretion ?
- No! There is no Problem.



## The Problem

● Mass Capture Rate

$$\dot{m}_c \sim \underline{10^{12} \text{ g/s}} \quad N_{\text{ISM}} V_6^{-3} M_{1.4}^2$$

● Luminosity

$$L_a \sim \underline{10^{32} \text{ erg/s}} \quad N_{\text{ISM}} V_6^{-3} M_{1.4}^3 R_6^{-1}$$

● Energy Range

$$\epsilon_\gamma \sim \underline{0.5 \text{ keV}} \quad L_{32}^{1/4} S_9^{-1/4}$$

## Predictions:

★ Treves & Colpi (1991, A&A, 241, 107)

$$\sim 5 \times 10^3$$

(ROSAT, All Sky Survey)

★ Blaes & Madau (1993, ApJ, 403, 690)

$$\sim 10^3 - 10^4$$

(ROSAT, All Sky Survey)

★ Popov et al. (2000, ApJ, 544, L53)

$$\sim 3 \times 10^4$$

Chandra & XMM-Newton

## Observations:

ROSAT, All Sky Survey

a few candidates...

- An accretion onto the surface of a NS can occur if

$$r_m = \left( \frac{\mu^2}{\dot{m}_c \sqrt{2GM}} \right)^{2/7} \lesssim r_{\text{cor}} = \left( \frac{GM}{\Omega^2} \right)^{1/3}$$

$$P_s \gtrsim P_0 \simeq \underline{7000 \text{ s}} \mu_{30}^{6/7} V_7^{9/7} N^{-3/7} M_{1.4}^{-11/7}$$

- The spin-down time scale of a NS is

$$\tau(P_0) = \tau_{\text{md}} + \tau_p \sim \underline{5 \times 10^9 \text{ yr}} \mu_{30}^{-1} V_7 N^{-1/2} I_{45} M_{1.4}^{-1}$$

$$\tau(P_0) \lesssim 10^{10} \text{ yr} \Rightarrow \underline{\underline{\text{NS is strongly magnetized}}} (r_m \gg r_{\text{ns}}) \text{ !!!}$$

# Geometry of the Accretion Flow

- Disk Accretion (in a turbulent ISM):

$$V_{\text{rel}} \approx \underline{10^6 \text{ cm/s}} \times \xi^{21/68} \mu_{30}^{-3/34} N^{3/68} M_{1.4}^{25/68} \left( \frac{V_t}{10^6 \text{ cm/s}} \right)^{21/68} \left( \frac{R_t}{10^{20} \text{ cm}} \right)^{-7/68}$$

- Spherical (Bondi) Accretion (Direct Accretion Approximation  $\dot{m}_c \equiv \dot{m}_a$ ):

$$V_{\text{rel}} < \underline{300 \text{ km/s}} \times N^{1/3} R_6^{-1/3} M_{1.4} \left( \frac{L_x}{10^{28} \text{ erg/s}} \right)^{-1/3}$$

- Kick velocity:  $V_{\text{rel}} \gtrsim \underline{5 \times 10^6 \text{ cm/s}}$

## Basic Conditions (normalization of parameters)

★ Strong magnetic field:  $\mu \simeq \underline{10^{30} \text{ G cm}^3}$

★ Relative velocity:  $50 \text{ km/s} \lesssim V_{\text{rel}} \lesssim 300 \text{ km/s}$

★ Spherical (Bondi) Accretion:  $\dot{m}_c \sim \underline{10^9 \text{ g/s}} \quad N \quad V_7^{-3} \quad M_{1.4}^2$

ESTIMATED NUMBER OF OBJECTS WITHIN 140 pc  $\sim \mathbf{10^4}$

# Flow Parameters at the magnetospheric boundary

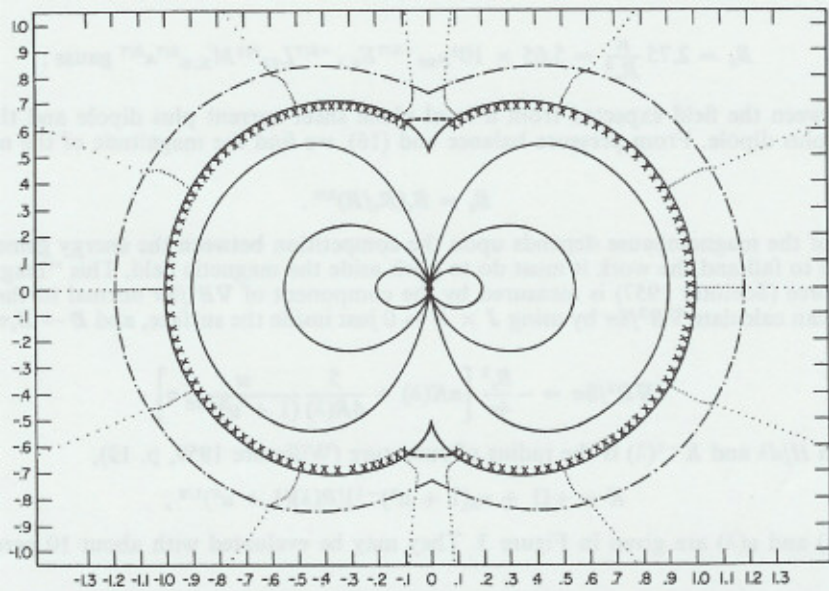
★ Magnetospheric Radius:  $r_m \simeq \underline{6 \times 10^{10} \text{ cm}}$   $\mu_{30}^{4/7} \dot{m}_9^{-2/7} M_{1.4}^{-1/7}$

★ Free-fall temperature:  $T_{\text{ff}} \simeq \underline{10^7 \text{ K}}$   $M_{1.4} \left( \frac{r_m}{6 \times 10^{10} \text{ cm}} \right)^{-1}$

★ Plasma density:  $N(r_m) \simeq \underline{300 \text{ cm}^{-3}}$   $\mu_{30}^2 T_7^{-1} \left( \frac{r_m}{6 \times 10^{10} \text{ cm}} \right)^{-6}$

★ Free-fall time scale:  $t_{\text{ff}} \simeq \underline{740 \text{ s}}$   $M_{1.4}^{-1/2} \left( \frac{r_m}{6 \times 10^{10} \text{ cm}} \right)^{3/2}$

★ Bremsstrahlung cooling time:  $t_{\text{br}} \sim \underline{10^5 \text{ yr}}$   $T_7^{1/2} \left( \frac{N(r_m)}{300 \text{ cm}^{-3}} \right)^{-1}$



# PLASMA ENTRY INTO THE MAGNETOSPHERE (I)

## Interchange instabilities

(Arons & Lea 1976, ApJ, 207, 914;      Elsner & Lamb 1976, Nature, 262, 356)

$$g_{\text{eff}} = \frac{GM_{\text{ns}}}{r_{\text{m}}^2(\lambda)} - \frac{V_{\text{T}_i}^2(r_{\text{m}})}{R_{\text{cur}}(\lambda)} > 0, \quad \longleftrightarrow \quad T(r_{\text{m}}) < 0.3T_{\text{ff}}(r_{\text{m}})$$

A Transient X-ray Source

Recurrent time:       $\sim \underline{10^5 \text{ yr}}$

Outburst duration:       $\sim \underline{15 \text{ min.}}$

Luminosity       $\sim 10^{29} \text{ erg/s}$

(Lamb, Fabian, Pringle, & Lamb 1977, ApJ, 217, 197)



# PLASMA ENTRY INTO THE MAGNETOSPHERE (II)

## Diffusion

★ Entry Rate:

$$\dot{M}_{\text{in}} \simeq \dot{M}_{\text{B}} = 2 \times 10^6 \text{ g s}^{-1} \alpha_{0.1}^{1/2} \mu_{30}^{-1/14} M_{1.4}^{11/7} N^{11/14} V_7^{-33/14}$$

★ Persistent Luminosity:

$$L_{\text{a}} \simeq \underline{2 \times 10^{26} \text{ erg s}^{-1}} \alpha_{0.1}^{1/2} \mu_{30}^{-1/14} N^{11/14} V_7^{-33/14} M_{1.4}^{19/7} r_6^{-1}$$

The Evolutionary Track of a Neutron Star  
contains **TWO** states of *Propeller*

(Davies, Fabian, & Pringle 1979, MNRAS, 186, 779)

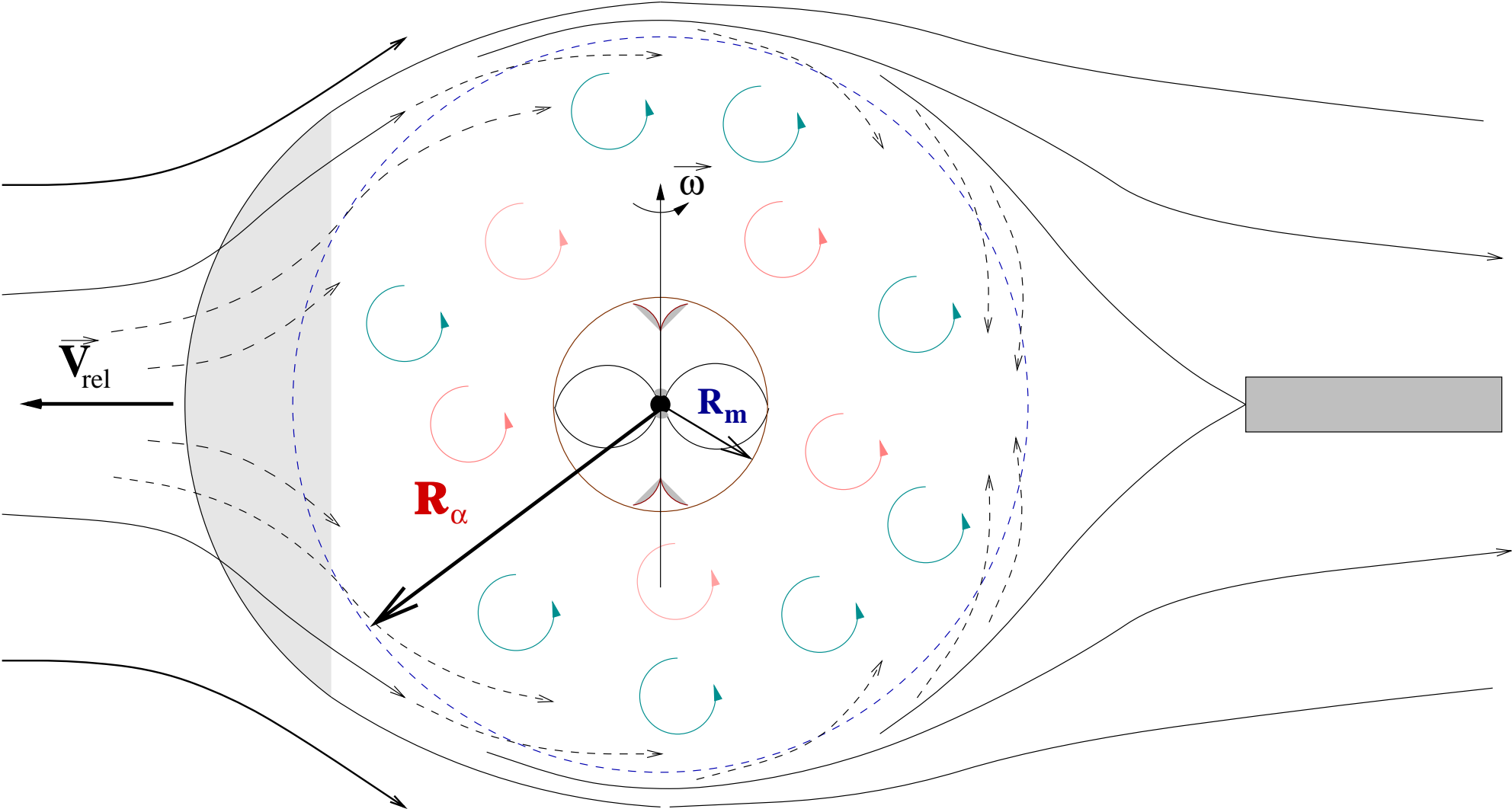
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I. EJECTOR (spin-powered pulsar)  $\implies$

II★. SUPERSONIC PROPELLER  $\implies$

III★★. SUBSONIC PROPELLER  $\implies$

IV. ACCRETOR (direct accretion:  $\dot{m}_c \equiv \dot{m}_a$ )



# Subsonic Propeller $\equiv$ Diffusion-driven Accretor

★ Spin-down heating:  $P_s \lesssim P_{\text{br}} \simeq \underline{10^5 \text{ s}} \quad \mu_{30}^{16/21} N^{-5/7} V_7^{15/7} M_{1.4}^{-34/21}$

$$\tau_{\text{br}} \sim \underline{2 \times 10^5 \text{ yr}} \quad \mu_{30}^{-2} I_{45} P_5 M_{1.4}$$

★ Heating by Radial Plasma Drift:

$$L_{\text{dr}} \sim \dot{m}_a \frac{GM_{\text{ns}}}{r}$$

Heating by the drift dominates cooling ( $t_{\text{heat}} \lesssim t_{\text{br}}$ ) if

$$\dot{m}_c \lesssim \dot{m}_0 \simeq 10^{14} \text{ g s}^{-1} \alpha_{0.1}^{7/17} \mu_{30}^{-1/17} V_7^{14/17} M_{1.4}^{16/17}$$

# Old Isolated NS Accreting Material from ISM can appear as

## ★ Persistent Sources

$$L_a \lesssim \underline{2 \times 10^{26} \text{ erg s}^{-1}} \alpha_{0.1}^{1/2} \mu_{30}^{-1/14} N^{11/14} V_7^{-33/14} M_{1.4}^{19/7} r_6^{-1}$$

$$P_s \gtrsim \underline{7000 \text{ s}} \mu_{30}^{6/7} V_7^{9/7} N^{-3/7} M_{1.4}^{-11/7}$$

## ★ Transient Sources

$$\underline{\text{Recurrent time:}} \sim \underline{10^5 \text{ yr}}$$

$$\underline{\text{Outburst duration:}} \sim \underline{15 \text{ min}}$$

$$\underline{\text{Luminosity}} \sim 10^{29} \text{ erg/s}$$