

Pulse Shape Modeling for accreting msec Pulsars

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Outline

- Motivation: accreting ms pulsars
- Properties: rapid rotation, general relativity
- Methods: exact and approximate calculations
- Some results: exact light curves; approximate light curve fits to observations

Motivation: accreting ms pulsars

Rotating neutron star with surface hot spot yields pulsations

Aim: to model the pulsations to obtain N^* parameters

Circular antipodal spots in Schwarzschild: Pechenick, Ftacelas & Cohen, 1983, ApJ, 274, 876

Main application has been to slowly rotating X-ray pulsars (Leahy 2004 ApJ 613, 517; Kraus et al 2003 ApJ 590, 424 and references therein)

Application to ms pulsars recent (Miller & Lamb 1998 ApJ 499 L37, Poutanen & Gierlinski 2003 MNRAS 343 1302, Cadeau Leahy Morsink 2005 ApJ 618 481)

Light curve generation

- Shape of emission region
- Surface emissivity (energy spectrum, angular distribution of emitted photons)
- Propagation of light to the observer
(Schwarzschild metric vs. exact neutron star metric)

Light propagation in Schwarzschild: Pechenick, Ftaclas & Cohen, 1983

Simple geodesics in Schwarzschild metric:

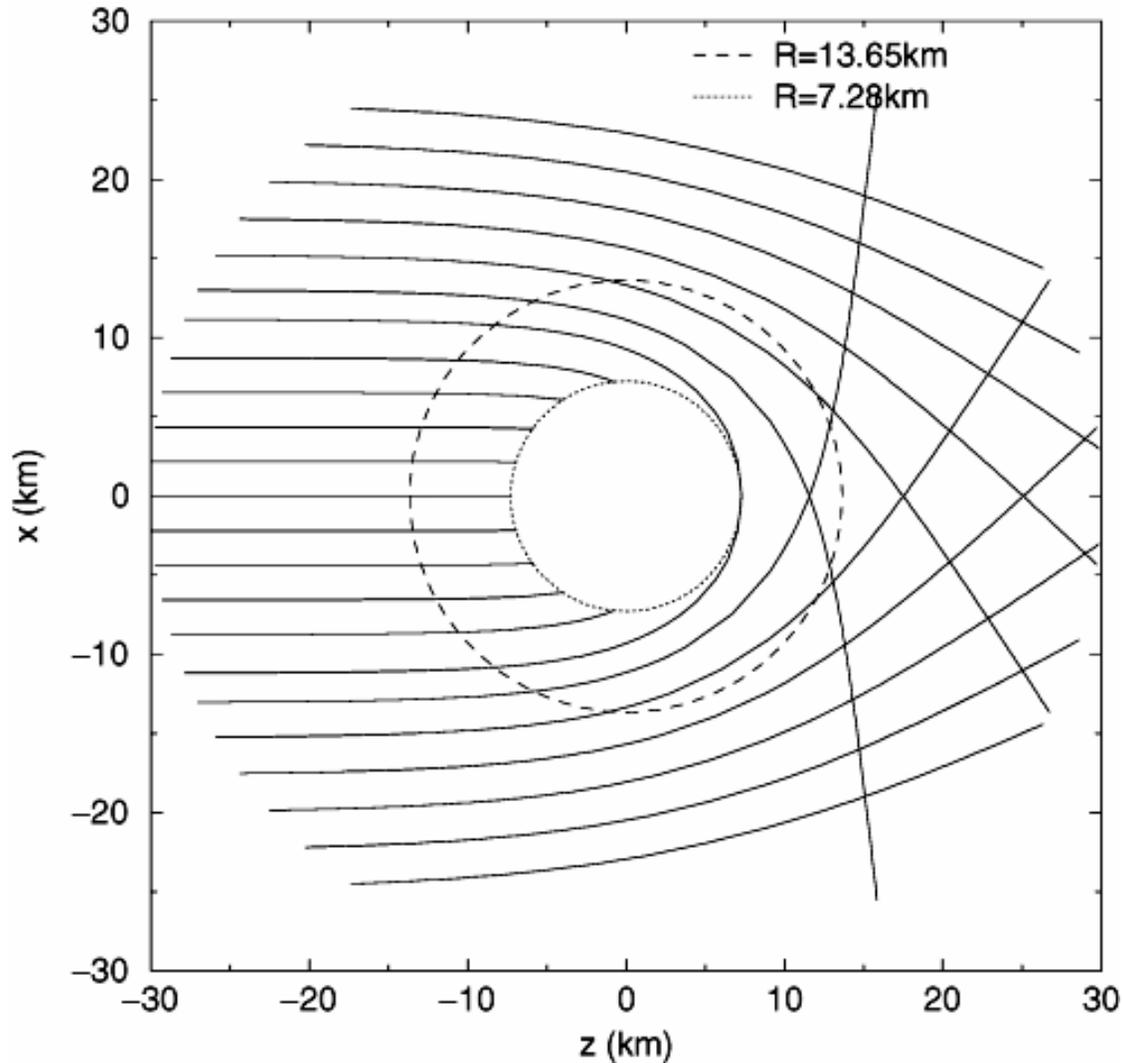
Let us now suppose that the photons are emitted from the surface ($r = R$) of an opaque sphere (the neutron star). The photon described by equation (2.2) obeys (Misner, Thorne, and Wheeler 1973, p. 673)

$$\left(\frac{1}{r^2} \frac{dr}{d\phi} \right)^2 + B^{-2}(r) = \frac{1}{b^2}, \quad (2.7)$$

where

$$B^{-2}(r) = (1 - 2M/r)/r^2. \quad (2.8)$$

Propagation of light



Ray paths for
Schwarzschild to
 $z \rightarrow -\infty$ (observer)

($R_{\text{sch}}=4.15\text{km}$)

Shadow zone

Part of back surface
is visible

Stellar surface is
magnified

Light bending for surface emission I

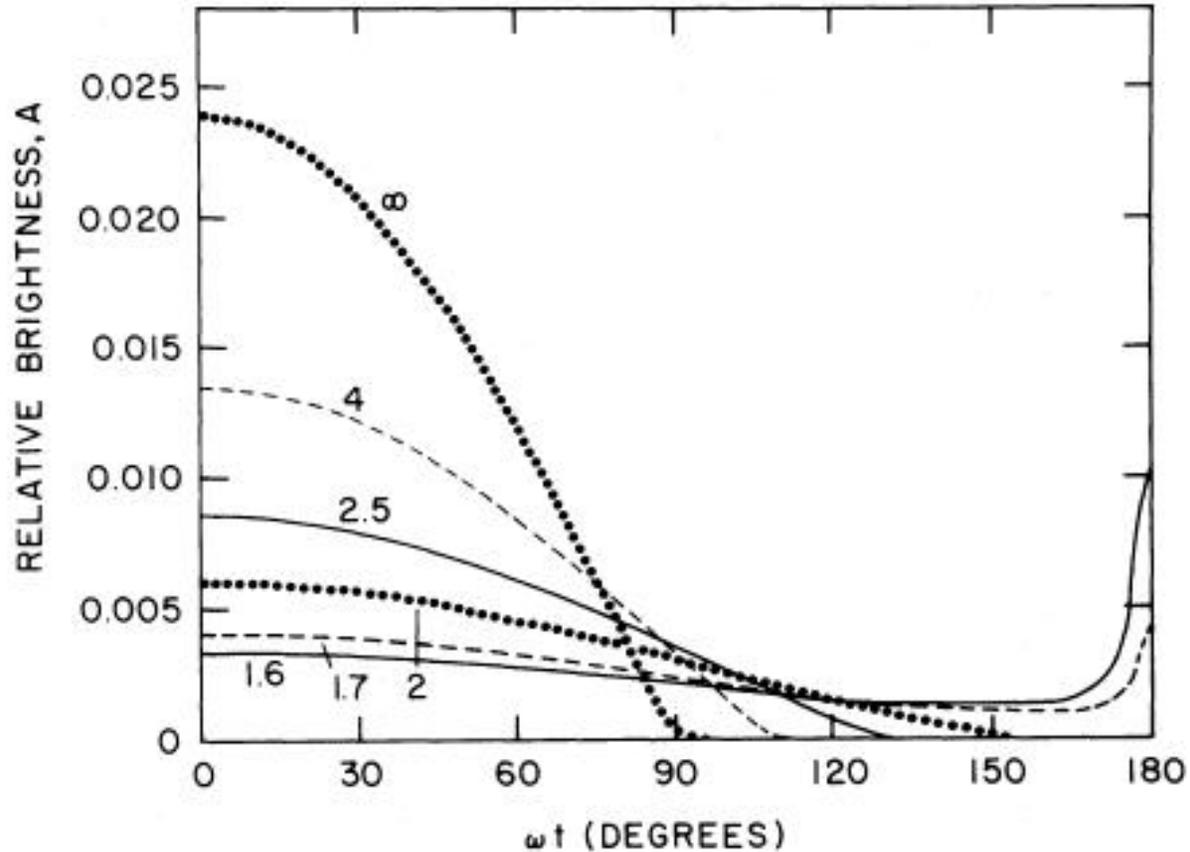
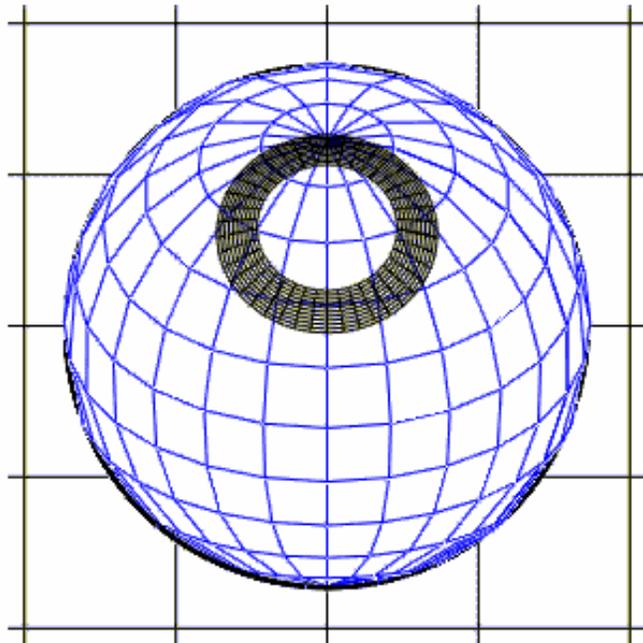


FIG. 4.—Relative brightness A (one polar cap) vs. ωt for $\alpha = 5^\circ$, $f(\delta) = 1$, $\beta = \gamma = 90^\circ$. Numbers denote values of $R/2M$.

Pechenick, Ftaclas & Cohen, 1983

Light bending for surface emission II

Neutron Star with Emission Regions (with Gravitational Light-bending)



More surface
visible

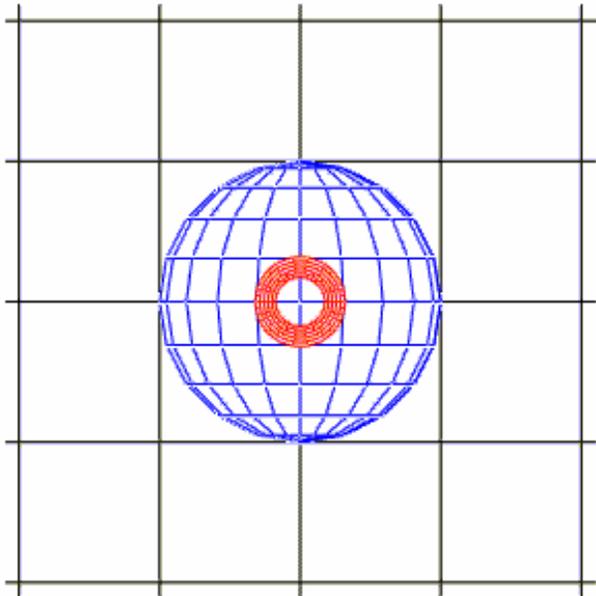
Magnification
increases toward
limb

Projected area of
emission surface
goes to zero at limb

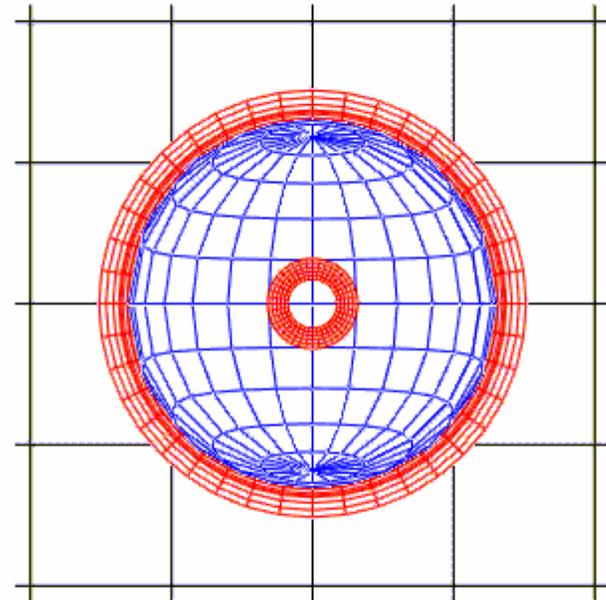
($R/R_{\text{sch}}=2.5$)

Light bending: column I ($R/R_s=2.5$)

Neutron Star with Emission Regions (Flat Space)



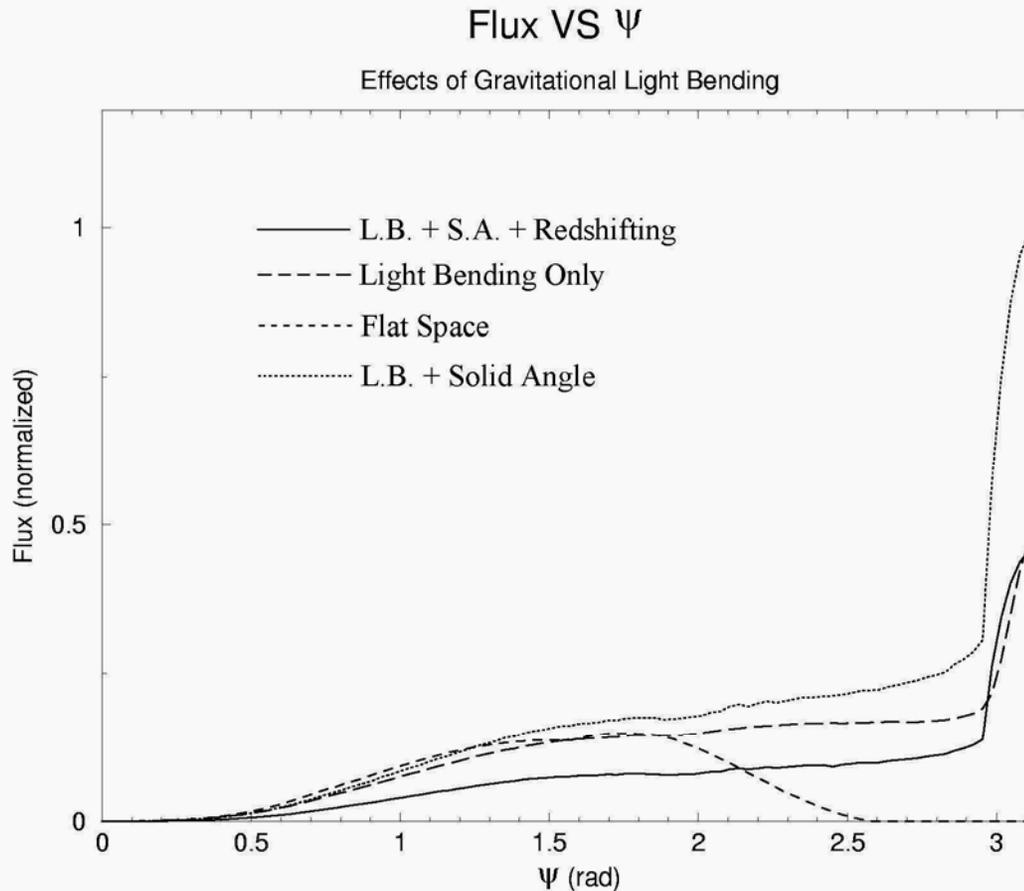
Neutron Star with Emission Regions (with Gravitational Light-bending)



$(x_1, y_1, z_1), (x_2, y_2, z_2)$

$(gx_0, gy_0, gz_0), (gx_1, gy_1, gz_1), (gx_2, gy_2, gz_2), (gx_3, gy_3, gz_3), (gx_4, gy_4, gz_4)$

Light bending: column II



Column exterior
emission only

“Spike” of
emission
due to
magnification
at limb

Area normal to
light
ray also max at
limb

What is so different for ms pulsars? I

Ms period means speeds at surface are not small compared to c ($v=47000\text{km/s}$ for $P=2\text{ms}$, $R=15\text{km}$)

Special-relativity:

- light aberration effects

- Doppler boosting and de-boosting

Time delays (R/c) are significant compared to the time for the star to rotate

What is so different for ms pulsars? II

Neutron star is oblate due to fast rotation

Spacetime metric is axisymmetric and must be computed numerically

We use algorithm of Cook, Shapiro and Teukolsky 1994; public domain code of Stergioulas and Friedman 1995

Equations of state from Arnett & Bowers 1977:

EOS A: very soft

EOS L: very stiff

Light propagation in exact metric:

Axisymmetric spacetime:

$$ds^2 = -e^{\gamma+\rho} dt^2 + e^{2\alpha} (d\bar{r}^2 + \bar{r}^2 d\theta^2) + e^{\gamma-\rho} \bar{r}^2 \sin^2 \theta (d\phi - \omega dt)^2$$

Geodesic equations:

$$\dot{t} = e^{-(\gamma+\rho)}(1 - \omega b)$$

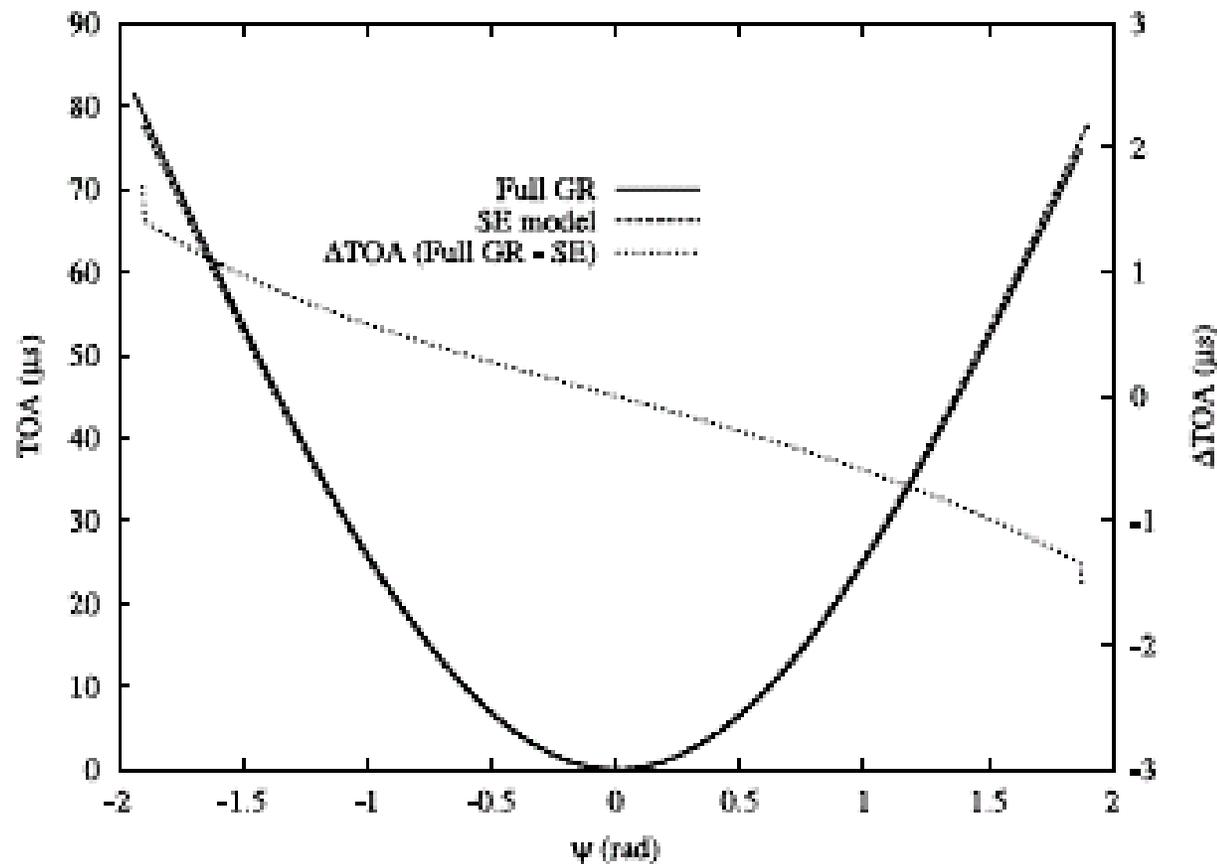
$$\dot{\phi} = \omega e^{-(\gamma+\rho)}(1 - \omega b) + e^{\rho-\gamma} \frac{b}{\bar{r}^2 \sin^2 \theta}$$

$$\ddot{\bar{r}} = -\alpha_{,\bar{r}} \left(\dot{\bar{r}}^2 - \bar{r}^2 \dot{\theta}^2 \right) - 2\alpha_{,\theta} \dot{\bar{r}} \dot{\theta} + \bar{r} \dot{\theta}^2 + \frac{1}{2} e^{-2\alpha} \mathcal{B}_{,\bar{r}}$$

$$\ddot{\theta} = \alpha_{,\theta} \left(\frac{\dot{\bar{r}}^2}{\bar{r}^2} - \dot{\theta}^2 \right) - 2 \left(\alpha_{,\bar{r}} + \frac{1}{\bar{r}} \right) \dot{\bar{r}} \dot{\theta} + \frac{1}{2\bar{r}^2} e^{-2\alpha} \mathcal{B}_{,\theta}$$

$$\mathcal{B} = e^{-(\gamma+\rho)}(1 - \omega b)^2 - \frac{b^2 e^{\rho-\gamma}}{\bar{r}^2 \sin^2 \theta}$$

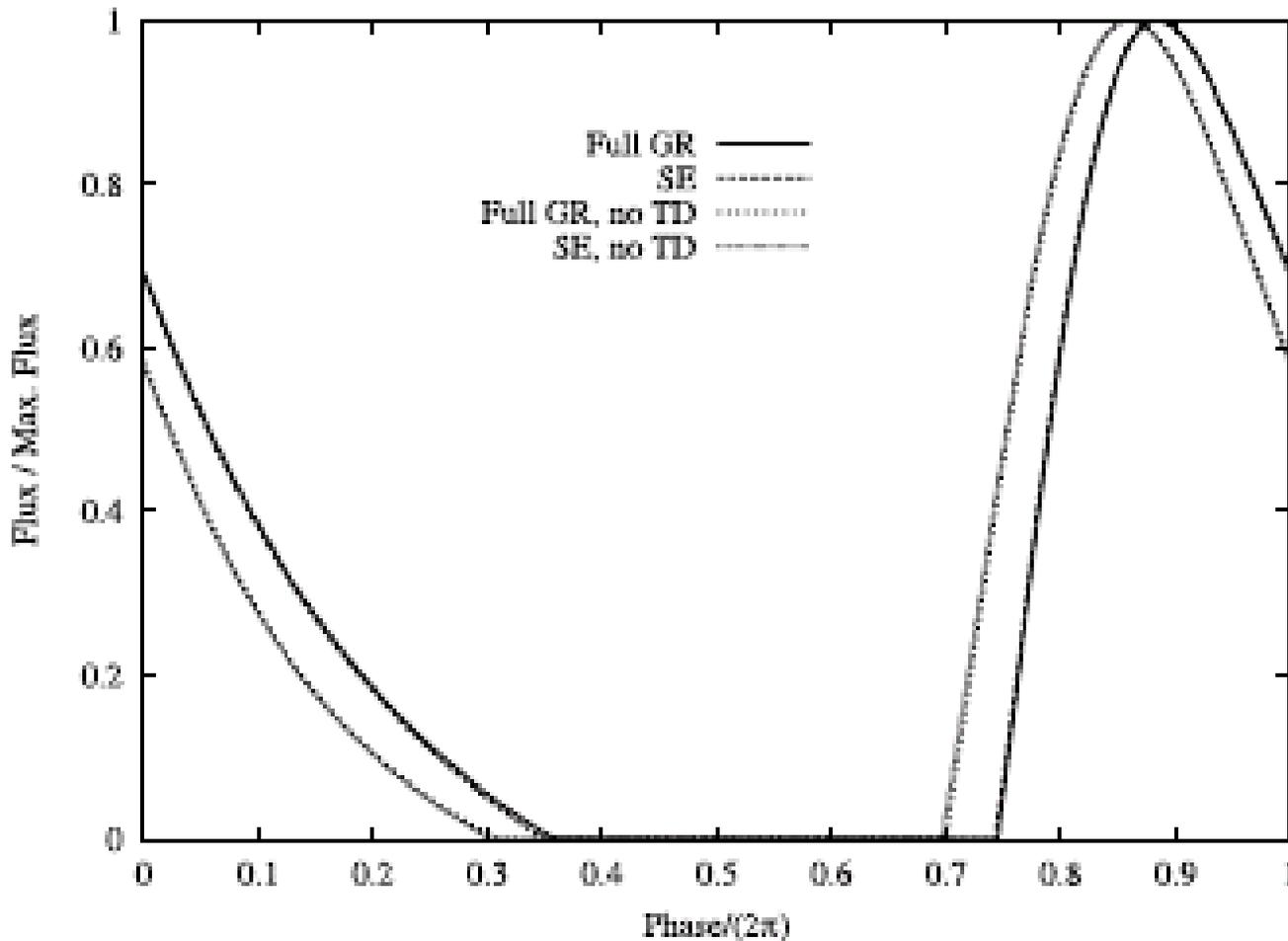
Time delays, equatorial rays



600 Hz
1.4 Msun
EOS L
 $R_{\text{eq}} = 16.38\text{km}$

FIG. 3.—TOA as a function of bending angle for model 4, calculated using the full metric and the SE metric.

Pulse shape, equatorial rays

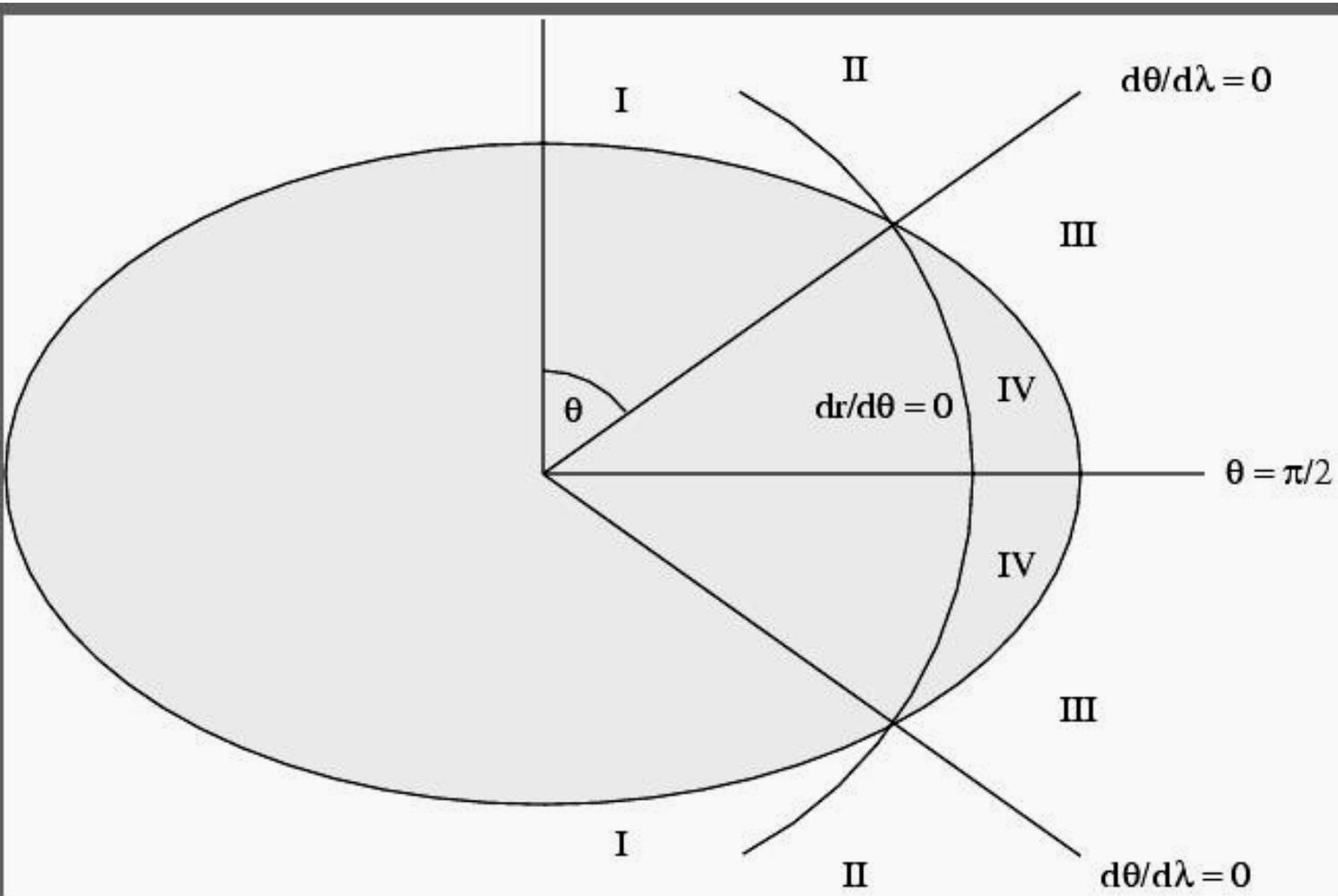


Asymmetric
pulse shape
due to:

1. Doppler
boosting
and
aberration
2. Time
delays

FIG. 5.—Bolometric pulse shapes for model 4 using different calculation methods.

Oblateness: visibility

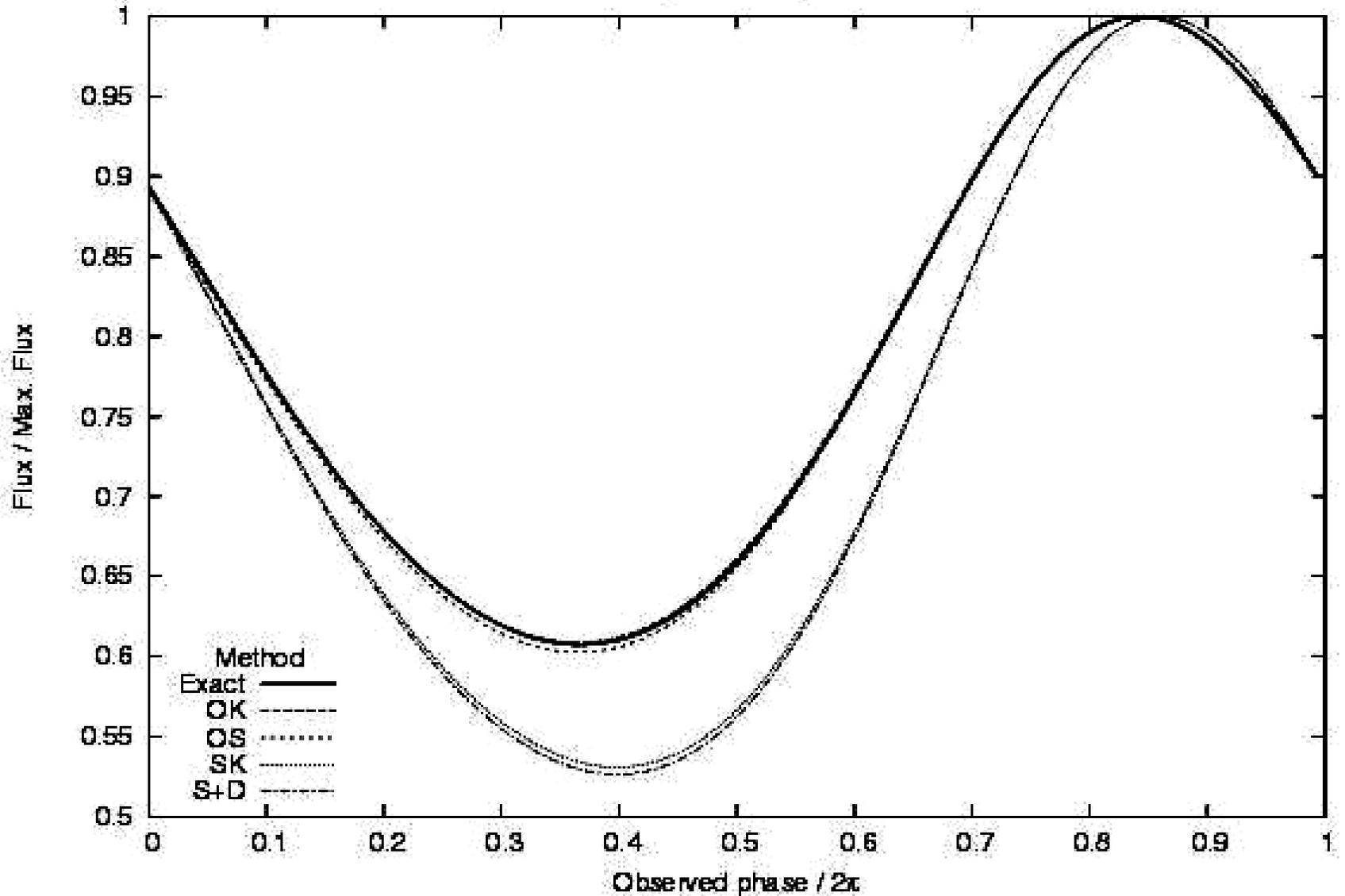


Spherical
star:
Rays
propagate
In regions
II,III,IV

Oblate star:
Rays
propagate
In regions
I,II,III

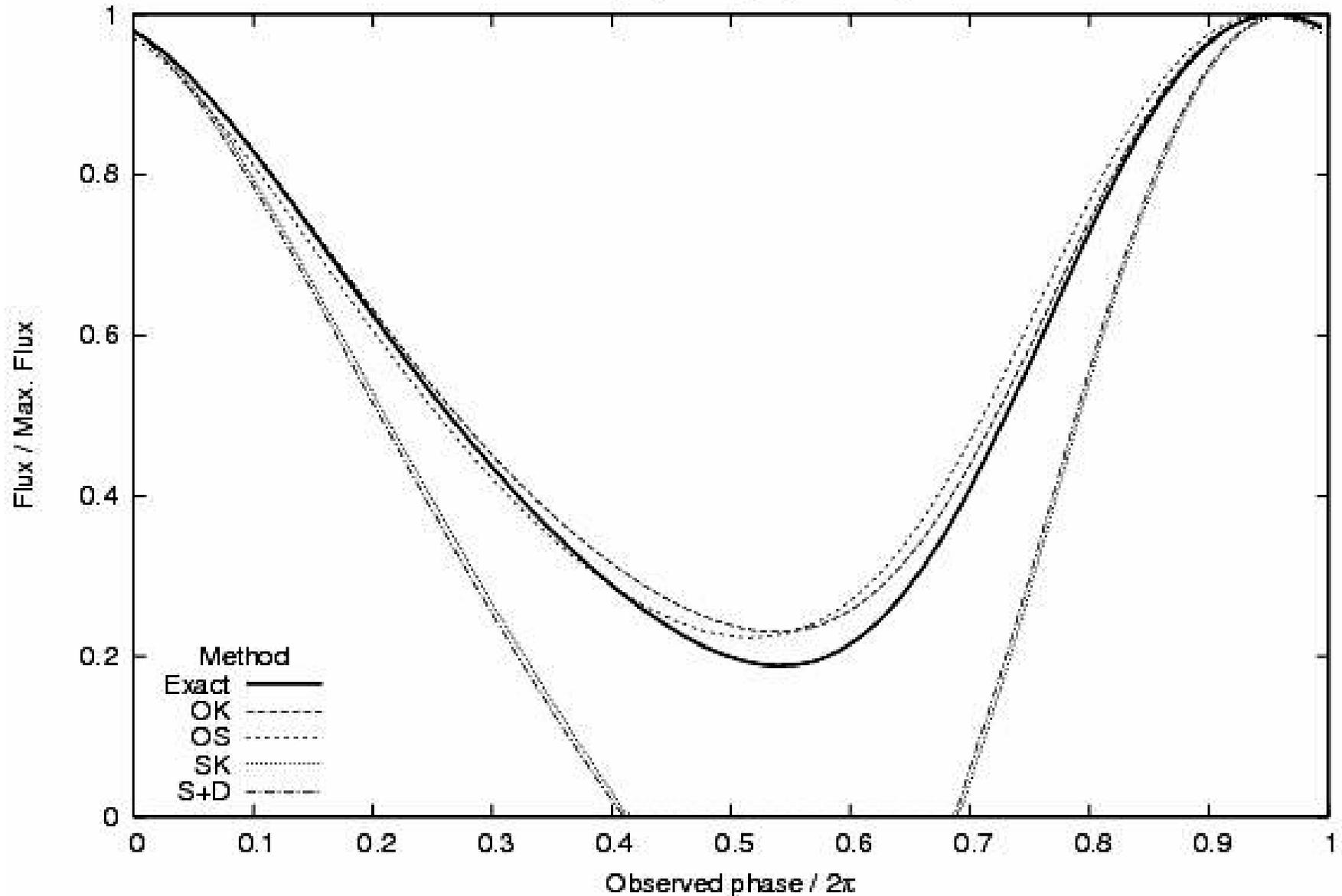
Different calculation methods I

L-600, $\theta_e=41$ deg, $\theta_o=20$ deg

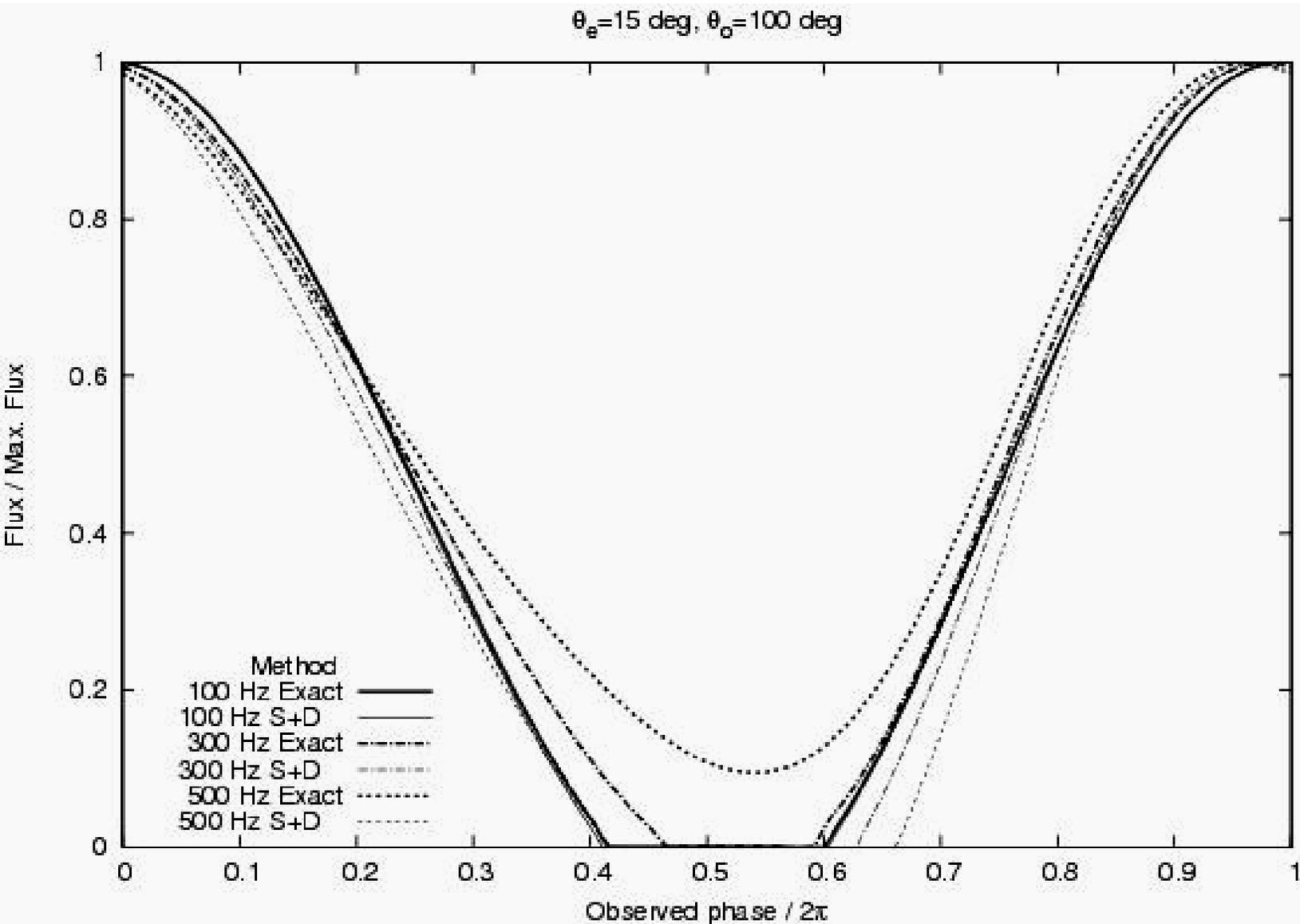


Different calculation methods II

L-600, $\theta_e=15$ deg, $\theta_o=100$ deg

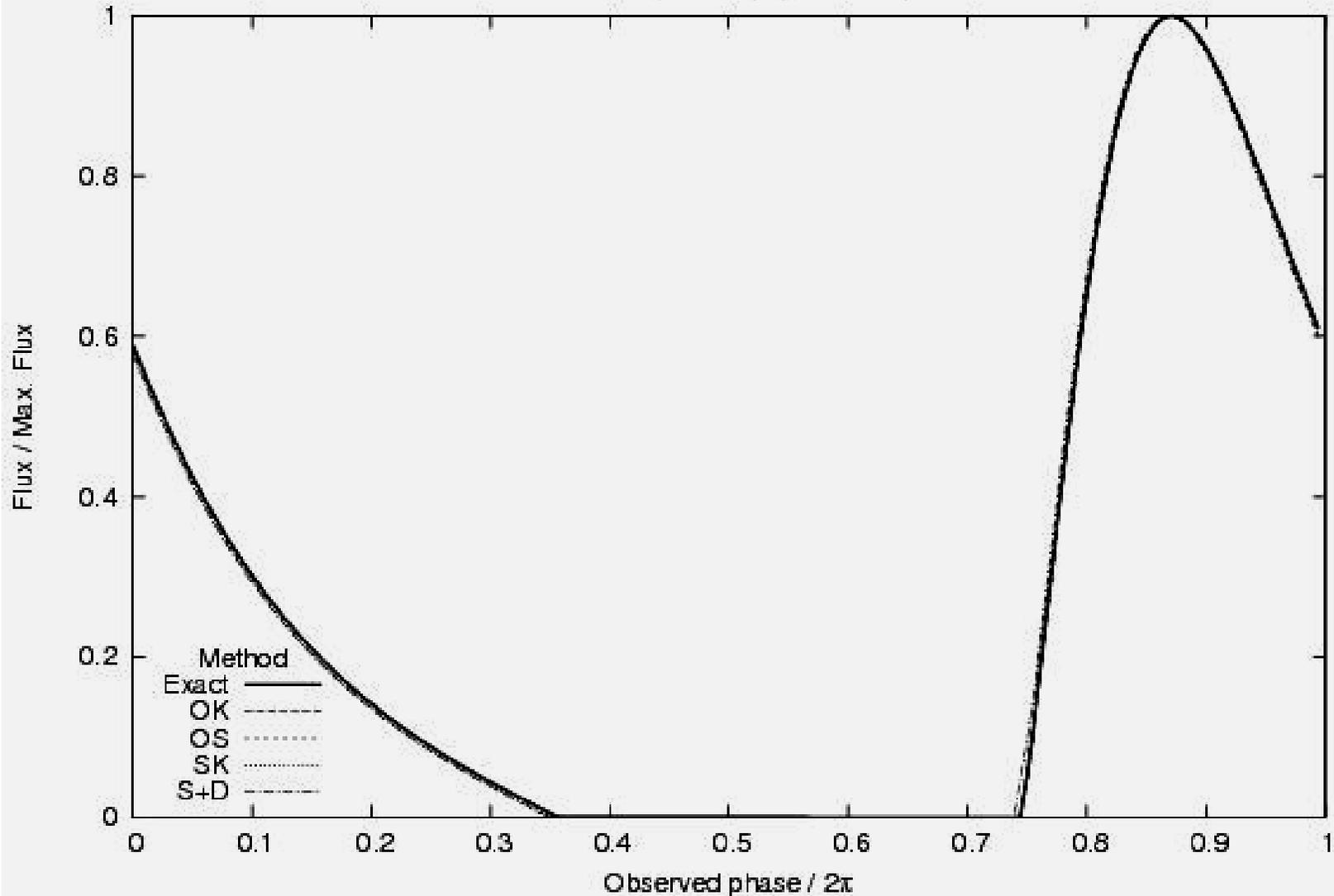


Different rotation rates



Near-equatorial rays

L-600, $\theta_e=85$ deg, $\theta_o=100$ deg



S+D model fit to exact light curves

Table 1. Comparison of True and Fitted Parameters of Neutron Stars. All neutron stars have a mass of $1.4M_{\odot}$.

θ_e	θ_o	EOS	Ω_{\star} (Hz)	M/M_{\odot}	$R(\theta_e)$ (km)		θ_e (deg)	θ_o (deg)	$GM/c^2R(\theta_e)$			χ^2
true				fit	true	fit	fit	fit	true	fit	unc.	
41°	100°	A	100	1.48	9.57	10.2	80.5	139.2	0.216	0.215	0.011	0.1
		L		1.32	14.83	13.8	80.5	133.0	0.140	0.142	0.027	0.02
		A	300	1.49	9.58	10.0	79.8	138.0	0.216	0.220	0.005	1
		L		1.09	14.82	11.1	67.0	95.6	0.140	0.145	0.024	0.3
		A	400	1.45	9.58	9.55	80.8	134.9	0.216	0.225	0.006	2
		L		1.17	14.80	11.9	58.0	96.3	0.140	0.145	0.023	0.4
		A	500	1.51	9.59	9.89	80.2	136.9	0.216	0.225	0.005	3
		L		1.29	14.78	12.7	52.7	98.1	0.140	0.15	0.02	0.8
		A	600	1.58	9.60	10.2	41.9	102.2	0.215	0.230	0.007	4
		L		1.30	14.74	12.0	57.9	97.5	0.140	0.160	0.015	2

S+D fits, cont'd

35°	100°	A	100	1.48	9.57	10.4	87.4	110.8	0.216	0.210	0.008	1
		L		1.45	14.86	15.3	84.0	103.7	0.139	0.140	0.027	0.05
		A	200	1.46	9.59	10.1	84.4	103.9	0.216	0.215	0.006	2
		L		1.43	14.95	15.6	86.7	107.6	0.138	0.135	0.029	0.4
		A	300	1.48	9.62	10.2	83.1	103.5	0.215	0.215	0.025	4
		L		1.70	15.10	17.9	80.0	123.3	0.137	0.140	0.015	0.4
		A	400	1.49	9.66	10.3	77.2	99.7	0.214	0.215	0.003	7
		L		1.40	15.35	16.0	85.0	105.5	0.135	0.130	0.009	4
		A	500	1.68	9.71	11.3	67.7	111.0	0.213	0.220	0.004	5
		L		1.51	15.73	17.9	61.7	97.7	0.131	0.125	0.009	6
		A	600	1.62	9.78	11.1	72.5	113.6	0.211	0.215	0.003	0.1
		L		1.40	16.35	17.3	77.8	102.6	0.127	0.120	0.007	0.2

Polar cap model for SAX J1808

Known pulse period
(401 Hz)

Pulse shape is nearly
symmetric

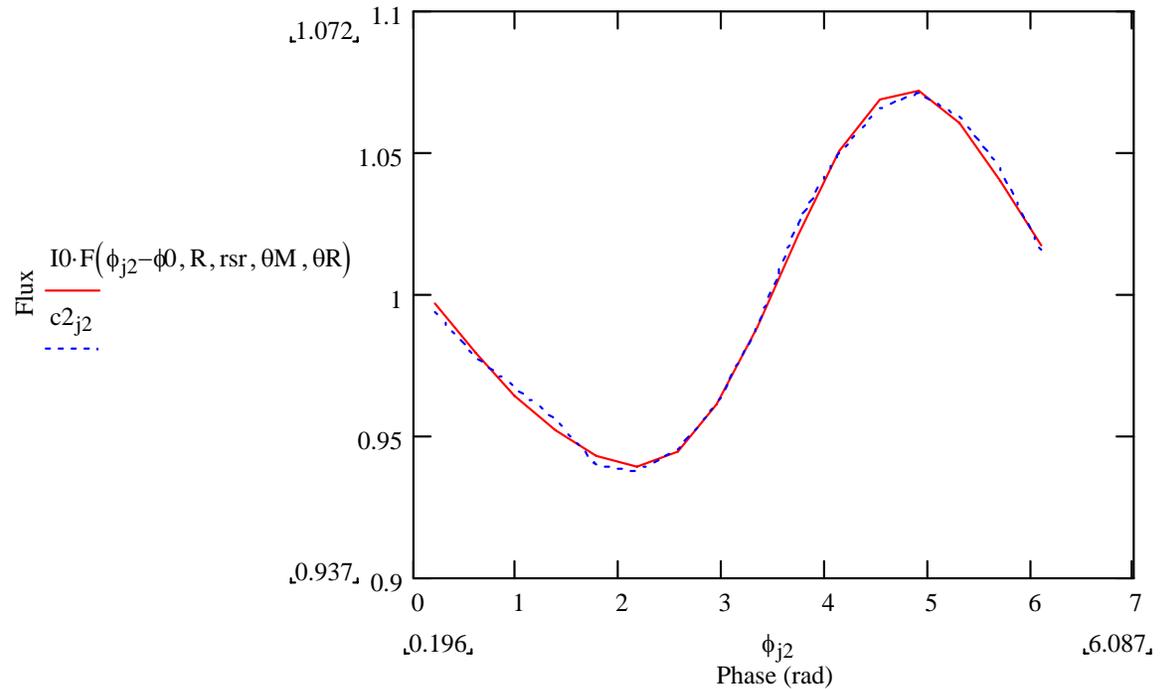
Constraints on model
parameters are loose

$R_s/R = .65$, $R = 8.1$ km

$M = 1.6$ Msun if no time
delays included

$R_s/R = .59$, $R = 6.3$ km

$M = 1.3$ Msun if time
delays included



12-18 keV

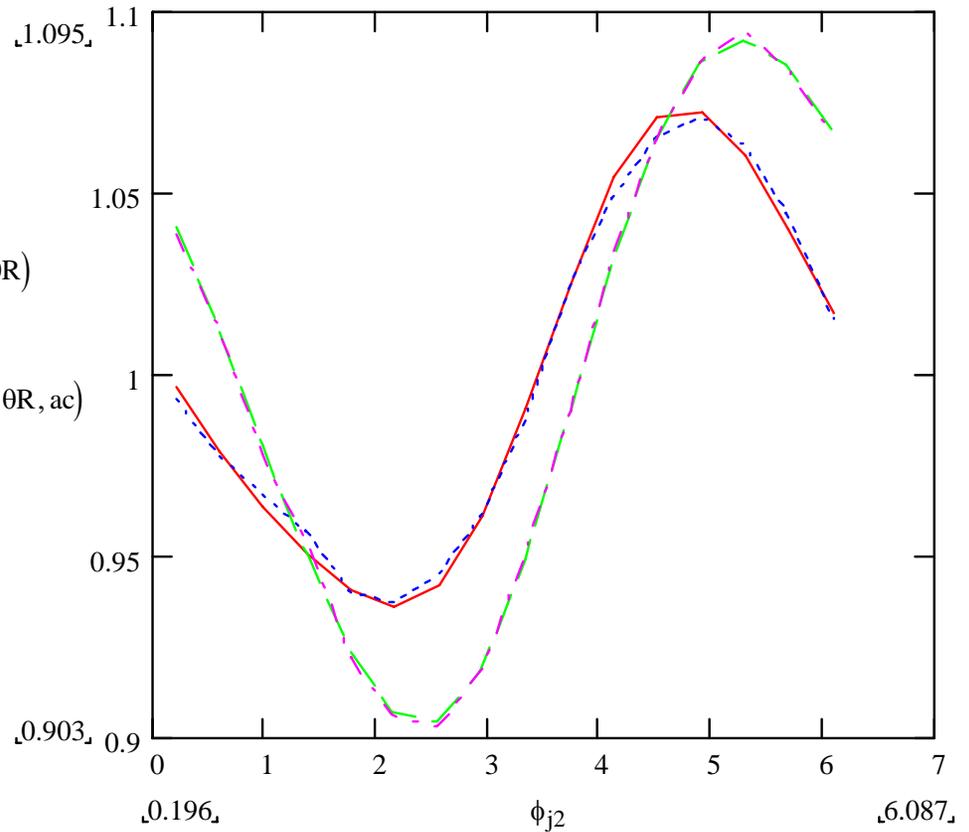
Two energy bands fit

Much better constrained

$R_s/R = .65$, $R = 7.2$ km
 $M = 1.6$ Msun if no time delays included

$R_s/R = .59$, $R = 8.3$ km
 $M = 1.65$ Msun if time delays included

- $I_0 \cdot F(\phi_{j2} - \phi_0, R, r_{sr}, \theta_M, \theta_R)$
- c_{2j2}
- $I_b \cdot F_2(\phi_{j2} - \phi_0, R, r_{sr}, \theta_M, \theta_R, ac)$
- c_{1j2}



3-4 keV: blackbody + Compton
 12-18 keV: Compton

New data Oct.2002

Best fit:

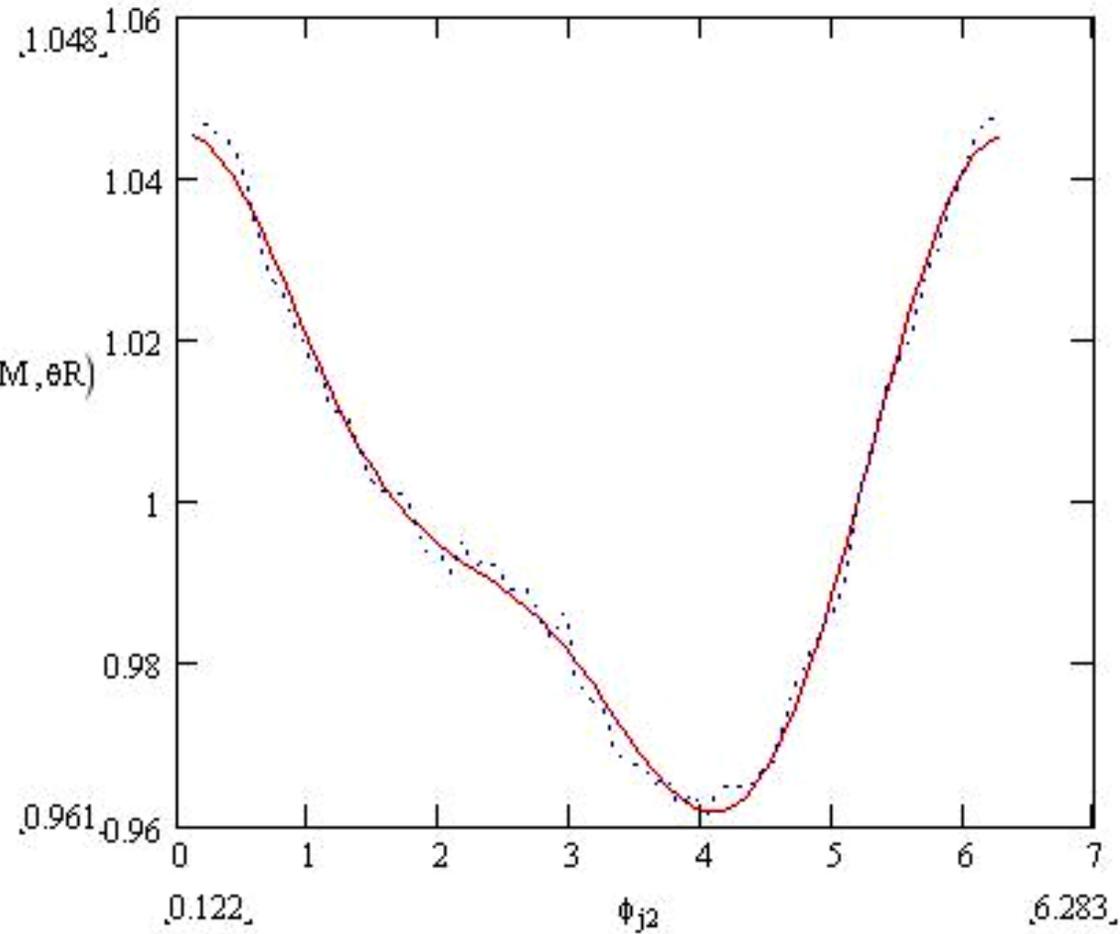
$R=4.8\text{km}$

$M=0.69M_{\text{sun}}$

$10 \cdot \text{Fit}(\phi_{j2}-\phi_0, R, r_{\text{sr}}, \theta M, \theta R)$

c^2_{j2}

.....



Joint fit: 1998 & 2002

R=6.0 km

M=0.94 Msun

R/Rs=2.15

Rotation axis

68deg

Hot spot:

17deg(Apr98)

12deg(Oct02)

Oblateness not
yet included

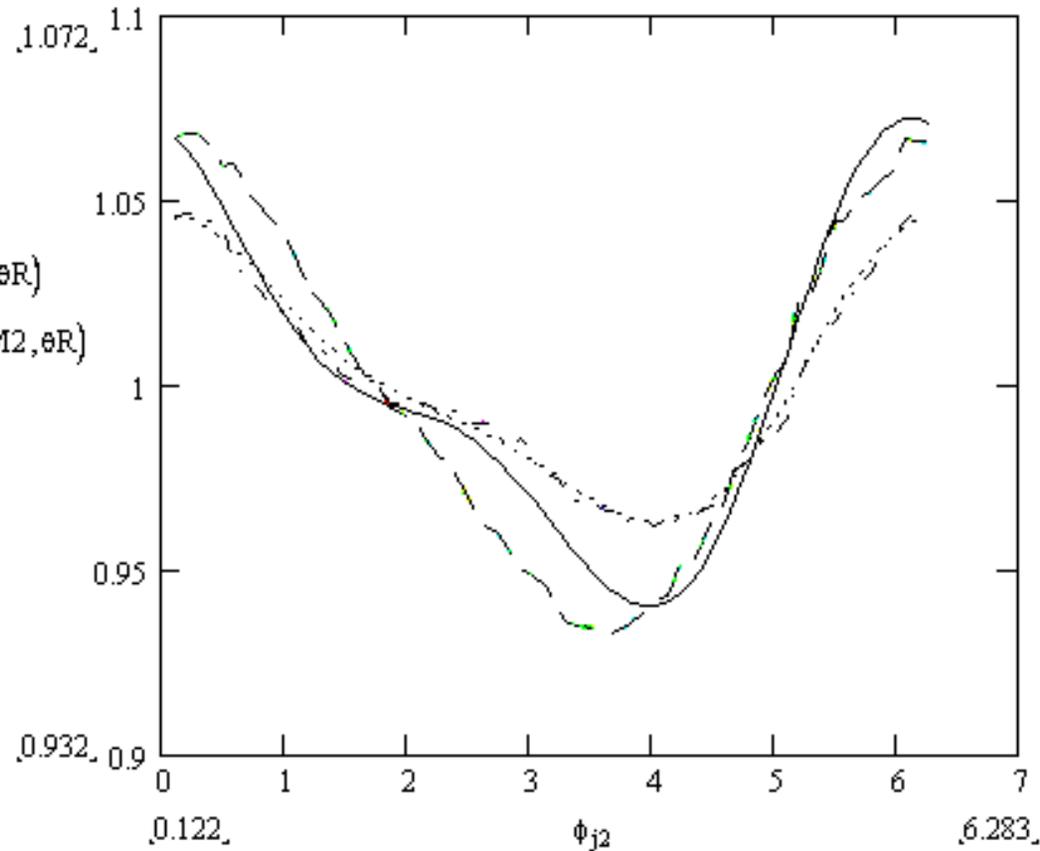
$I0 \cdot Ft(\phi_{j2} - \phi0, R, rsr, \theta M, \theta R)$

$I02 \cdot Ft(\phi_{j2} - \phi02, R, rsr, \theta M2, \theta R)$

.....
c1j2

— · —
c2j2

— · —
.....



Fit results II

$R=6.0$ km, $M=0.94$ Msun: Quark star??

Not yet.

Compare to fitting to exact calc for small angles: M_{fit} & R_{fit} can be too large or too small by up to 50% due to missing oblateness

More realistic guess: $R < 10$ km, $M < 1.5$ Msun

Coming soon: approx. model with oblateness.

Summary

- New aspects of exact calculations: ray tracing in full metric, time delays, oblate neutron star
- Oblateness and time delays are largest effects
- Exact calc. too slow for multi-parameter fitting to data
- Approx. calc: currently incorporating oblateness
- Fitting results sensitive to including all effects
- Goal: obtain M/R , M , R and emission region parameters for ms pulsars

