Pulsar Winds

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The Crab Nebula Central star is source of particles and magnetic field (Piddington 1957) and waves (Rees & Gunn 1974).

- Few particles: magnetic dipole radiation? Damping \Rightarrow propagation only for $\omega_{pe} < \Omega$. For Crab: $r > 10^8 r_L$ (e.g., Melatos & Melrose 1996)
- Many particles, MHD wind + shock

- The wind–nebula connection
 - MHD simulation
- Acceleration of the wind
 - Dissipation in shocks/current sheets
- Observation of the wind
 - Optical pulse shapes and polarisation
 - TeV emission binary system PSR B1259 –63
- Acceleration of particles
 - Two mechanisms?



2D relativistic MHD by at least three groups

Komissarov & Lyubarsky 2003; Khangoulian & Bogovalov 2003; Del Zanna et al 2004

Key ingredients:

- relativistic, anisotropic wind (power $\propto \sin^2 \theta$)
- low magnetisation σ (at least near equator)

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- Radial flow, toroidal magnetic field
- Jet formed downstream of termination shock
- Low value of $\langle \sigma \rangle \sim 0.03$
- No constraints on μ parameter (= $L/\dot{M}c^2$) from the dynamics
- Problems with the inner ring (knots, front/back brightness ratio) — may reflect kinetic effects, such as proposed for the wisps

(Gallant & Arons 1994; Spitkovsky & Arons 2000)

Exact solution for force-free, split monopole (Michel 1973): no collimation, $B_{\phi} \propto \sin \theta / r$ (no closed field lines)

Super-(magneto)sonic flow: $\Gamma \rightarrow \text{constant}$ (Bogovalov 1997)

$$\sigma = \frac{B^2/8\pi}{\Gamma nmc^2}$$
$$= \text{constant}$$

the σ problem

Accelerate the wind:

- Collimation? Not for monopole-like flows (e.g., Bogovalov & Tsinganos 1999) but in principle possible (Vlahakis 2004)
- Dissipation? Oblique rotator (Coroniti 1990) and damping of wave component — how fast?

Problem not really a problem:

- σ still high after the shock (Begelman 1998)?
 Difficult to recover nice pictures...
- the (striped) field dissipates in the termination shock (Lyubarsky 2003) Transition must remain thin

Acceleration of the wind

Dissipation forced by charge starvation $(B \propto 1/r, n \propto 1/r^2)$

Entropy wave or FMS wave (small wavelength approx. $r \gg r_{\rm L}$) Lyubarsky & Kirk 2001; Lyubarsky 2003;

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Current sheets



Magnetic pressure balanced by hot plasma in sheet. Key question: What controls the dissipation rate?

Short wavelength approximation (Kirk & Skjæraasen 2003)

Slow dissipation	Tearing-mode	Fast
Coroniti (1980);	Lyubarsky (1996)	Drenkhahn & Spruit (2002)
Michel (1994);		
Lyubarsky & Kirk (2001)		
$\Gamma \propto r^{1/2}$	$\Gamma \propto r^{5/12}$	$\Gamma \propto r^{1/3}$
$\frac{r_{\max}}{L} = \hat{L}^{1/2}$	$\frac{r_{\rm max}}{2} = \mu^{4/5} \hat{L}^{3/10}$	$\frac{r_{\text{max}}}{2} = \mu^{4/5} \hat{L}^{3/10}$
$r_{ m L}$ $-$	$r_{ m L}$ $r_{ m L}$ –	$r_{\rm L}$ $r_{\rm L}$ –

 $\hat{L} = L(\pi^2 e^2/m^2 c^5)$, (= 1.5×10^{22} for Crab)

No consistent conversion mechanism for $\mu > 10\hat{L}^{1/4}$

- Gamma-rays from unshocked wind
 - Targets from companion star: swamped by emission from shocked wind (Ball & Kirk 2000)
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TeV emission from shocked wind in PSR B1259 –63

- Hadronic emission (Kawachi et al 2004)
- Inverse Compton model

Unique pulsar/Be star binary



$$\frac{r_{\rm Be}}{r_{\rm p}} = \sqrt{\frac{L_{\rm s.d.}}{\dot{M}v_{\rm wind}c}} \sim 0.7$$

But $L_{\rm s.d.} \gg L_{\rm Be wind}$



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Observations with H.E.S.S. telescopes

Schlenker et al (2005), Aharonian et al astro-ph/0506280

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 Connect with acceleration theory?

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Crab-type injection $dn/d\gamma \propto \gamma^{-1.6}$ at $\gamma < \Gamma_{\rm peak}$

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